A martingale approach for the elephant random walk

Bernard Bercu®

Institut de Mathématiques de Bordeaux, Université de Bordeaux, UMR 5251, 351 Cours de la Libération, 33405 Talence cedex, France

E-mail: bernard.bercu@u-bordeaux.fr

Received 29 August 2017 Accepted for publication 24 October 2017 Published 29 November 2017



Abstract

The purpose of this paper is to establish, via a martingale approach, some refinements on the asymptotic behavior of the one-dimensional elephant random walk (ERW). The asymptotic behavior of the ERW mainly depends on a memory parameter p which lies between zero and one. This behavior is totally different in the diffusive regime $0 \le p < 3/4$, the critical regime p = 3/4, and the superdiffusive regime 3/4 . In the diffusive and critical regimes, we establish some new results on the almost sure asymptotic behavior of the ERW, such as the quadratic strong law and the law of the iterated logarithm. In the superdiffusive regime, we provide the first rigorous mathematical proof that the limiting distribution of the ERW is not Gaussian.

Keywords: elephant random walk, martingales, strong law of large numbers, asymptotic normality

1. Introduction

Random walks with long memory arose naturally in applied mathematics, theoretical physics, computer sciences and econometrics. One of them is the so-called elephant random walk (ERW). It is a one-dimensional discrete-time random walk on integers, which has a complete memory of its whole history. It was introduced in 2004 by Schütz and Trimper [18] in order to investigate the long-term memory effects in non-Markovian random walks. It was referred to as the ERW in allusion to the famous saying that elephants can remember where they have been

A wide range of literature is available on the asymptotic behavior of the ERW and its extensions [1, 2, 5, 6, 7, 8, 11, 14]. However, many things remain to be done. The goal of this paper is to answer to several natural questions on the ERW.

One of them concerns the influence of the memory parameter *p* on the almost sure asymptotic behavior of the ERW. Quite recently, Baur and Bertoin [1] and independently Coletti *et al*

[5] studied the asymptotic normality of the ERW in the diffusive regime p < 3/4 as well as in the critical regime p = 3/4. However, very few results are available on the almost sure asymptotic behavior of the ERW in the regime $p \le 3/4$. We shall fill the gap by proving a quadratic strong law as well as a law of the iterated logarithm for the ERW in the regime $p \le 3/4$.

Another key question concerns the limiting distribution of the ERW in the superdiffusive regime p > 3/4. Initially, it was suggested by Schütz and Trimper [18] that, even in the superdiffusive regime, the ERW has a Gaussian limiting distribution. Later, it was conjectured by Da Silva *et al* [8] that this limiting distribution is not Gaussian, see also [5, 16]. One can observe that the analytical study of [8] is not sufficient to prove the non-gaussianity of the limiting distribution. We shall provide a rigorous mathematical proof that the limiting distribution is not Gaussian. Starting from the symmetric initial condition, we will also show that this limiting distribution is sub-Gaussian.

Baur and Bertoin [1] extensively used the connection to Pólya-type urns [17] as well as two functional limit theorems for multitype branching processes due to Janson [10], see also [4]. Our strategy is totally different as it relies on the theory of martingales. To be more precise, we shall make use of the strong law of large numbers and the central limit theorem for martingales [9, 12] as well as the law of the iterated logarithm for martingales [19, 20]. We strongly believe that our approach could be successfully extended to ERWs with stops [6, 13], to amnesiac ERWs [7], as well as to multi-dimensional ERWs [6, 15].

The paper is organized as follows. In section 2, we introduce the exact ERW and the martingale we shall extensively make use of. Section 3 is devoted to the main results of the paper. We establish the almost sure asymptotic behavior as well as the asymptotic normality of the ERW in the diffusive and critical regimes. Moreover, in the superdiffusive regime, we provide the first rigorous mathematical proof that the limiting distribution of the ERW is not Gaussian. Our martingale approach is described in appendix A, while all technical proofs of section 3 are postponed to appendices B and C.

2. The elephant random walk

The one-dimensional ERW is defined as follows. The random walk starts at the origin at time zero, $S_0 = 0$. At time n = 1, the elephant moves to the right with probability q and to the left with probability 1 - q where q lies between zero and one. Hence, the position of the elephant at time n = 1 is given by $S_1 = X_1$ where X_1 has a Rademacher $\mathcal{R}(q)$ distribution. Afterwards, at any time $n \ge 1$, we choose uniformly at random an integer k among the previous times $1, \ldots, n$, and we define

$$X_{n+1} = \begin{cases} +X_k & \text{with probability} & p, \\ -X_k & \text{with probability} & 1-p, \end{cases}$$

where the parameter $p \in [0, 1]$ is the memory of the ERW. Then, the position of the ERW is given by

$$S_{n+1} = S_n + X_{n+1}. (2.1)$$

In order to understand well how the elephant moves, it is straightforward to see that for any time $n \ge 1$, $X_{n+1} = \alpha_n X_{\beta_n}$ where α_n and β_n are two independent discrete random variables, where α_n has a Rademacher $\mathcal{R}(p)$ distribution while β_n is uniformly distributed over the integers $\{1, \dots, n\}$. Moreover, α_n is independent of X_1, \dots, X_n .

Let (\mathcal{F}_n) be the increasing sequence of σ -algebras, $\mathcal{F}_n = \sigma(X_1, \dots, X_n)$. For any time $n \ge 1$, we clearly have

$$\mathbb{E}[X_{n+1}|\mathcal{F}_n] = \mathbb{E}[\alpha_n] \times \mathbb{E}[X_{\beta_n}|\mathcal{F}_n] = (2p-1)\frac{S_n}{n} \text{ a.s.}$$
 (2.2)

Which, together with (2.1), implies that

$$\mathbb{E}[S_{n+1}|\mathcal{F}_n] = \gamma_n S_n \text{ where } \gamma_n = \left(\frac{n+2p-1}{n}\right). \tag{2.3}$$

Moreover,

$$\prod_{k=1}^{n} \gamma_k = \frac{\Gamma(n+2p)}{\Gamma(n+1)\Gamma(2p)}$$

where Γ stands for the Euler gamma function. Therefore, let (M_n) be the sequence of random variables defined, for all $n \ge 0$, by $M_n = a_n S_n$ where $a_1 = 1$ and, for all $n \ge 2$,

$$a_n = \prod_{k=1}^{n-1} \gamma_k^{-1} = \frac{\Gamma(n)\Gamma(2p)}{\Gamma(n+2p-1)}.$$
 (2.4)

Since $a_n = \gamma_n a_{n+1}$, we clearly deduce from (2.3) that for any time $n \ge 1$,

$$\mathbb{E}[M_{n+1}|\mathcal{F}_n]=M_n$$
 a.s.

In other words, the sequence (M_n) is a multiplicative real martingale. Our strategy is to make use of the martingale (M_n) in order to deduce the asymptotic behavior of (S_n) .

3. Main results

3.1. The diffusive regime

Our first result concerns the almost sure convergence of the ERW in the diffusive regime where $0 \le p < 3/4$.

Theorem 3.1. We have the almost sure convergence

$$\lim_{n \to \infty} \frac{S_n}{n} = 0 \text{ a.s.} \tag{3.1}$$

We now focus our attention on the almost sure rates of convergence of the ERW.

Theorem 3.2. We have the quadratic strong law

$$\lim_{n \to \infty} \frac{1}{\log n} \sum_{k=1}^{n} \left(\frac{S_k}{k}\right)^2 = \frac{1}{3 - 4p} \text{ a.s.}$$
(3.2)

In addition, we also have the law of the iterated logarithm

$$\limsup_{n \to \infty} \left(\frac{1}{2n \log \log n} \right)^{1/2} S_n = -\liminf_{n \to \infty} \left(\frac{1}{2n \log \log n} \right)^{1/2} S_n$$

$$= \frac{1}{\sqrt{3 - 4p}} \text{ a.s.}$$
(3.3)

In particular,

$$\limsup_{n \to \infty} \frac{S_n^2}{2n \log \log n} = \frac{1}{3 - 4p} \text{ a.s.}$$
(3.4)

Our next result is devoted to the asymptotic normality of the ERW in the diffusive regime $0 \le p < 3/4$.

Theorem 3.3. We have the asymptotic normality

$$\frac{S_n}{\sqrt{n}} \xrightarrow{\mathcal{L}} \mathcal{N}\left(0, \frac{1}{3 - 4p}\right). \tag{3.5}$$

3.2. The critical regime

Hereafter, we investigate the critical regime where the memory parameter p = 3/4.

Theorem 3.4. We have the almost sure convergence

$$\lim_{n \to \infty} \frac{S_n}{\sqrt{n} \log n} = 0 \text{ a.s.}$$
 (3.6)

The almost sure rates of convergence of the ERW are as follows.

Theorem 3.5. We have the quadratic strong law

$$\lim_{n \to \infty} \frac{1}{\log \log n} \sum_{k=2}^{n} \left(\frac{S_k}{k \log k}\right)^2 = 1 \text{ a.s.}$$
(3.7)

In addition, we also have the law of the iterated logarithm

$$\limsup_{n \to \infty} \left(\frac{1}{2n \log n \log \log \log n} \right)^{1/2} S_n = -\liminf_{n \to \infty} \left(\frac{1}{2n \log n \log \log \log n} \right)^{1/2} S_n$$

$$= 1 \text{ a.s.}$$
(3.8)

In particular,

$$\limsup_{n \to \infty} \frac{S_n^2}{2n \log n \log \log \log n} = 1 \text{ a.s.}$$
(3.9)

One can observe a very unusual rate of convergence in the law of the iterated logarithm. Our next result deals with the asymptotic normality of the ERW in the critical regime p = 3/4.

Theorem 3.6. We have the asymptotic normality

$$\frac{S_n}{\sqrt{n\log n}} \xrightarrow{\mathcal{L}} \mathcal{N}(0,1). \tag{3.10}$$

3.3. The superdiffusive regime

Finally, we focus our attention on the more complicated superdiffusive regime where 3/4 .

Theorem 3.7. We have the almost sure convergence

$$\lim_{n \to \infty} \frac{S_n}{n^{2p-1}} = L \text{ a.s.}$$

$$\tag{3.11}$$

where L is a non-degenrate random variable. This convergence also holds in \mathbb{L}^4 , which means that

$$\lim_{n \to \infty} \mathbb{E}\left[\left|\frac{S_n}{n^{2p-1}} - L\right|^4\right] = 0. \tag{3.12}$$

Remark 3.1. One can observe that the first three moments of S_n were previously calculated in [8] in the special case q = 1. However, the analytical study of [8] is not sufficient to evaluate the moments of L.

Theorem 3.8. The first four moments of L are given by

$$\mathbb{E}[L] = \frac{2q-1}{\Gamma(2p)},\tag{3.13}$$

$$\mathbb{E}[L^2] = \frac{1}{(4p-3)\Gamma(2(2p-1))},\tag{3.14}$$

$$\mathbb{E}[L^3] = \frac{2p(2q-1)}{(2p-1)(4p-3)\Gamma(3(2p-1))},\tag{3.15}$$

$$\mathbb{E}[L^4] = \frac{6(8p^2 - 4p - 1)}{(8p - 5)(4p - 3)^2\Gamma(4(2p - 1))}.$$
(3.16)

Remark 3.2. Our last result provides the first rigorous mathematical proof that, in the superdiffusive regime, the limiting distribution L of the ERW is not Gaussian. We denote by μ and σ^2 the mean value and the variance of L, $\mu = \mathbb{E}[L] = (2q-1)/\Gamma(2p)$ and $\sigma^2 = \mathbb{E}[(L-\mu)^2]$. Moreover, let α and κ be the skewness and the kurtosis of L, respectively defined by

$$\alpha = \frac{\mathbb{E}[(L-\mu)^3]}{\sigma^3}$$
 and $\kappa = \frac{\mathbb{E}[(L-\mu)^4]}{\sigma^4}$.

In the special case q = 1/2, the skewness $\alpha = 0$, while the kurtosis

$$\kappa = \frac{6(8p^2 - 4p - 1)(\Gamma(2(2p - 1)))^2}{(8p - 5)\Gamma(4(2p - 1))}.$$

It is not hard to see that κ is a decreasing function of p such that, for all 3/4 . This means that <math>L has a sub-Gaussian distribution. Furthermore, if $q=1, \alpha>0$ and if $q=0, \alpha<0$. Finally, the kurtosis κ shares the same value for q=1 or q=0, which can be smaller, greater or equal to 3.

Acknowledgment

The author would like to thank Adrien Richou who makes him discover the elephant kingdom.

Appendix A. The martingale approach

We already saw that the sequence (M_n) given, for all $n \ge 0$, by $M_n = a_n S_n$, is a multiplicative real martingale. Moreover, for any $n \ge 1$, X_n is a binary random variable taking values in $\{+1, -1\}$. Consequently, $|S_n| \le n$, which implies that (M_n) is a locally square integrable. The martingale (M_n) can be rewritten in the additive form

$$M_n = \sum_{k=1}^n a_k \varepsilon_k \tag{A.1}$$

since its increments $\Delta M_n = M_n - M_{n-1}$ satisfy $\Delta M_n = a_n S_n - a_{n-1} S_{n-1} = a_n \varepsilon_n$ where $\varepsilon_n = S_n - \gamma_{n-1} S_{n-1}$. The predictable quadratic variation [9] associated with (M_n) is given by $\langle M \rangle_0 = 0$ and, for all $n \ge 1$,

$$\langle M \rangle_n = \sum_{k=1}^n \mathbb{E}[\Delta M_k^2 | \mathcal{F}_{k-1}]. \tag{A.2}$$

We immediately obtain from (2.3) that $\mathbb{E}[\varepsilon_{n+1}|\mathcal{F}_n] = 0$. Moreover, it follows from (2.1) together with (2.2) that

$$\mathbb{E}[S_{n+1}^2|\mathcal{F}_n] = \mathbb{E}[S_n^2 + 2S_nX_{n+1} + 1|\mathcal{F}_n] = 1 + (2\gamma_n - 1)S_n^2 \text{ a.s.}$$
(A.3)

Consequently, as $\mathbb{E}[\varepsilon_{n+1}^2|\mathcal{F}_n] = \mathbb{E}[S_{n+1}^2|\mathcal{F}_n] - \gamma_n^2 S_n^2$, we deduce from (A.3) that, for all $n \ge 1$,

$$\mathbb{E}[\varepsilon_{n+1}^2 | \mathcal{F}_n] = 1 + (2\gamma_n - 1)S_n^2 - \gamma_n^2 S_n^2 = 1 - (\gamma_n - 1)^2 S_n^2 \text{ a.s.}$$

$$= 1 - (2p - 1)^2 \left(\frac{S_n}{n}\right)^2 \text{ a.s.}$$
(A.4)

By the same token,

$$\mathbb{E}[\varepsilon_{n+1}^4 | \mathcal{F}_n] = 1 - 3(\gamma_n - 1)^4 S_n^4 + 2(\gamma_n - 1)^2 S_n^2 \text{ a.s.}$$

$$= 1 - 3(2p - 1)^4 \left(\frac{S_n}{n}\right)^4 + 2(2p - 1)^2 \left(\frac{S_n}{n}\right)^2 \text{ a.s.}$$
(A.5)

On the one hand, if p=1/2, $\mathbb{E}[\varepsilon_{n+1}^2|\mathcal{F}_n]=1$ then $\mathbb{E}[\varepsilon_{n+1}^4|\mathcal{F}_n]=1$ a.s. On the other hand, we obtain from (A.4) and (A.5) the almost sure upper bounds

$$\sup_{n\geq 0} \mathbb{E}[\varepsilon_{n+1}^2 | \mathcal{F}_n] \leqslant 1 \text{ and } \sup_{n\geq 0} \mathbb{E}[\varepsilon_{n+1}^4 | \mathcal{F}_n] \leqslant \frac{4}{3}. \tag{A.6}$$

Hereafter, we deduce from (A.1), (A.2) and (A.4) that

$$\langle M \rangle_n = \sum_{k=1}^n a_k^2 - (2p-1)^2 \zeta_n \text{ where } \zeta_n = \sum_{k=1}^{n-1} a_{k+1}^2 \left(\frac{S_k}{k}\right)^2.$$
 (A.7)

The asymptotic behavior of the martingale (M_n) is closely related to the one of

$$v_n = \sum_{k=1}^n a_k^2 = \sum_{k=1}^n \left(\frac{\Gamma(k)\Gamma(2p)}{\Gamma(k+2p-1)} \right)^2.$$

Via standard results on the asymptotic behavior of the Euler gamma function, we have three regimes. In the diffusive regime where $0 \le p < 3/4$,

$$\lim_{n \to \infty} \frac{\nu_n}{n^{3-4p}} = \ell \text{ where } \ell = \frac{(\Gamma(2p))^2}{3-4p}.$$
 (A.8)

In the critical regime where p = 3/4,

$$\lim_{n \to \infty} \frac{v_n}{\log n} = \frac{\pi}{4}.\tag{A.9}$$

In the superdiffusive regime where $3/4 , <math>v_n$ converges to the finite value

$$\lim_{n \to \infty} v_n = \sum_{k=0}^{\infty} \left(\frac{\Gamma(k+1)\Gamma(2p)}{\Gamma(k+2p)} \right)^2 = \sum_{k=0}^{\infty} \frac{(1)_k (1)_k (1)_k}{(2p)_k (2p)_k k!} = {}_{3}F_{2} {1, 1, 1 \choose 2p, 2p} 1$$
(A.10)

where, for any $a \in \mathbb{R}$, $(a)_k = a(a+1)\cdots(a+k-1)$ for $k \ge 1$, $(a)_0 = 1$ stands for the Pochhammer symbol and ${}_3F_2$ is the generalized hypergeometric function defined by

$$_{3}F_{2}\binom{a,b,c}{d,e}|z\rangle = \sum_{k=0}^{\infty} \frac{(a)_{k}(b)_{k}(c)_{k}}{(d)_{k}(e)_{k}k!}z^{k}.$$

Appendix B. Proofs of the almost sure convergence results

B.1. The diffusive regime

Proof of theorem 3.1. First of all, we focus our attention on the proof of the almost sure convergence (3.1). We already saw from (A.7) that $\langle M \rangle_n \leq v_n$. Moreover, it follows from (A.8) that, in the diffusive regime, v_n increases to infinity with an arithmetic speed n^{3-4p} . Then, we obtain from the strong law of large numbers for martingales given, e.g. by theorem 1.3.24 of [9] that

$$\lim_{n\to\infty}\frac{M_n}{v_n}=0 \text{ a.s.}$$

Which implies that

$$\lim_{n \to \infty} \frac{M_n}{n^{3-4p}} = 0 \text{ a.s.}$$
 (B.1)

This convergence is not sharp enough to prove (3.1). However, thanks to the last part of theorem 1.3.24, we also have the almost sure rate of convergence in (B.1)

$$\frac{M_n^2}{v_n} = O(\log v_n) \text{ a.s.}$$

Which ensures that

$$\frac{M_n^2}{n^{3-4p}} = O(\log n) \text{ a.s.}$$
(B.2)

Hereafter, as $M_n = a_n S_n$, it clearly follows from (B.2) that

$$\frac{S_n^2}{n} = O(\log n) \text{ a.s.}$$

Which immediately leads to (3.1).

Proof of theorem 3.2. Denote by f_n the explosion coefficient associated with the martingale (M_n) given, for all $n \ge 1$, by

$$f_n = \frac{a_n^2}{v_n}.$$

We clearly obtain from (A.8) that f_n converges to zero. Moreover, it follows from the almost sure convergence (3.1) together with (A.4) that

$$\lim_{n\to\infty} \mathbb{E}[\varepsilon_{n+1}^2|\mathcal{F}_n] = 1 \text{ a.s.}$$

Furthermore, we already saw in (A.6) that

$$\sup_{n\geq 0} \mathbb{E}[\varepsilon_{n+1}^4 | \mathcal{F}_n] \leqslant \frac{4}{3} \text{ a.s.}$$

Consequently, by virtue of the quadratic strong law for martingales given, e.g., in theorem 3 of [3],

$$\lim_{n\to\infty}\frac{1}{\log \nu_n}\sum_{k=1}^n f_k\Big(\frac{M_k^2}{\nu_k}\Big)=1 \text{ a.s.}$$

Which implies, via (A.8), that

$$\lim_{n \to \infty} \frac{1}{\log n} \sum_{k=1}^{n} \frac{a_k^2 M_k^2}{v_k^2} = (3 - 4p) \text{ a.s.}$$
(B.3)

Therefore, as $M_n = a_n S_n$ and $n^2 a_n^4$ is equivalent to $(3 - 4p)^2 v_n^2$, we find from (B.3) that

$$\lim_{n \to \infty} \frac{1}{\log n} \sum_{k=1}^{n} \left(\frac{S_k}{k}\right)^2 = \frac{1}{3 - 4p} \text{ a.s.}$$
 (B.4)

Which completes the proof the quadratic strong law (3.2). We shall now proceed to the proof of the law of the iterated logarithm given by (3.3). In order to apply the law of the iterated logarithm for martingales due to Stout [20], see also corollary 6.4.25 in [9], it is only necessary to verify that

$$\sum_{n=1}^{+\infty} \frac{a_n^4}{v_n^2} < +\infty. \tag{B.5}$$

This is clearly satisfied since $a_n^4 v_n^{-2}$ is equivalent to $(3-4p)^2 n^{-2}$ and, as is well known,

$$\sum_{n=1}^{+\infty} \frac{1}{n^2} = \frac{\pi^2}{6}.$$

Hence, we find from the law of the iterated logarithm for martingales that

$$\limsup_{n \to \infty} \left(\frac{1}{2\nu_n \log \log \nu_n} \right)^{1/2} M_n = -\liminf_{n \to \infty} \left(\frac{1}{2\nu_n \log \log \nu_n} \right)^{1/2} M_n$$

$$= 1 \text{ a.s.}$$
(B.6)

As previously seen, the identity $M_n = a_n S_n$ together with (B.6) immediately lead to

$$\begin{split} \limsup_{n \to \infty} \left(\frac{1}{2n \log \log n} \right)^{1/2} S_n &= - \liminf_{n \to \infty} \left(\frac{1}{2n \log \log n} \right)^{1/2} S_n \\ &= \frac{1}{\sqrt{3 - 4p}} \text{ a.s.} \end{split}$$

Which completes the proof of theorem 3.2.

B.2. The critical regime

Proof of theorem 3.4. We shall proceed as in the proof of theorem 3.1. It follows from (A.9) that, in the critical regime where p = 3/4, v_n increases slowly to infinity with a logarithmic speed $\log n$. We obtain from theorem 1.3.24 of [9] that

$$\frac{M_n^2}{\log n} = O(\log \log n) \text{ a.s.}$$
(B.7)

However, we deduce from (2.4) that

$$\lim_{n \to \infty} na_n^2 = \frac{\pi}{4}.\tag{B.8}$$

Hence, as $M_n = a_n S_n$, we find from (B.7), together with (B.8), that

$$\frac{S_n^2}{n\log n} = O(\log\log n) \text{ a.s.}$$

Which immediately implies (3.6).

Proof of theorem 3.5. The proof of theorem 3.5 is left to the reader as it follows essentially the same lines as the one of theorem 3.2.

B.3. The superdiffusive regime

Proof of theorem 3.7. In the superdiffusive regime $3/4 , we already saw from (A.10) that <math>v_n$ converges to a finite value. Hence, as $\langle M \rangle_n \le v_n$, we deduce from (A.1), together with theorem 1.3.15 of [9], the almost sure convergence

$$\lim_{n \to \infty} M_n = M \text{ where } M = \sum_{k=1}^{\infty} a_k \varepsilon_k.$$
 (B.9)

Moreover, we obtain from (2.4) that

$$\lim_{n \to \infty} n^{2p-1} a_n = \Gamma(2p). \tag{B.10}$$

Consequently, as $M_n = a_n S_n$, (3.11) immediately follows from (B.9) and (B.10). One can observe that

$$L = \frac{1}{\Gamma(2p)} \sum_{k=1}^{\infty} a_k \varepsilon_k.$$

It only remains to prove convergence (3.12). For that purpose, it is only necessary to show that the martingale (M_n) is bounded in \mathbb{L}^4 . It is not hard to see that (M_n) is bounded in \mathbb{L}^2 . As a matter of fact, since $M_{n+1}^2 = (a_{n+1}\varepsilon_{n+1} + M_n)^2$, it follows from (A.6) that

$$\mathbb{E}[M_{n+1}^2|\mathcal{F}_n] = a_{n+1}^2 \mathbb{E}[\varepsilon_{n+1}^2|\mathcal{F}_n] + M_n^2 \leqslant a_{n+1}^2 + M_n^2 \text{ a.s.}$$

Taking expectation on both sides, we get that $\mathbb{E}[M_{n+1}^2] \leq a_{n+1}^2 + \mathbb{E}[M_n^2]$ leading to

$$\mathbb{E}[M_n^2] \leqslant \sum_{k=1}^n a_k^2. \tag{B.11}$$

Consequently, we obtain from (A.10) together with (B.11) that

$$\sup_{n \ge 1} \mathbb{E}[M_n^2] \le {}_3F_2 \binom{1, 1, 1}{2p, 2p} | 1$$
 $< \infty.$ (B.12)

By the same token, as $M_{n+1}^4 = (a_{n+1}\varepsilon_{n+1} + M_n)^4$, we clearly have

$$\mathbb{E}[M_{n+1}^4|\mathcal{F}_n] = \sum_{\ell=0}^4 \binom{4}{\ell} a_{n+1}^\ell M_n^{4-\ell} \mathbb{E}[\varepsilon_{n+1}^\ell|\mathcal{F}_n]. \tag{B.13}$$

On the one hand, we already saw from (A.6) that

$$\mathbb{E}[\varepsilon_{n+1}^2|\mathcal{F}_n] \leqslant 1$$
 and $\mathbb{E}[\varepsilon_{n+1}^4|\mathcal{F}_n] \leqslant \frac{4}{3}$ a.s.

On the other hand, as in (A.4), we also have

$$\mathbb{E}[\varepsilon_{n+1}^3 | \mathcal{F}_n] = 2(\gamma_n - 1)^3 S_n^3 - 2(\gamma_n - 1)^2 S_n \text{ a.s.}$$

$$= 2(2p - 1) \frac{S_n}{n} \left((2p - 1)^2 \left(\frac{S_n}{n} \right)^2 - 1 \right) \text{ a.s.}$$
(B.14)

Therefore, as $M_n = a_n S_n$ and $|S_n| \leq n$, we obtain from (B.14) that

$$M_n\mathbb{E}[\varepsilon_{n+1}^3|\mathcal{F}_n]\leqslant 0$$
 a.s.

Consequently, it follows from (B.13) that

$$\mathbb{E}[M_{n+1}^4|\mathcal{F}_n] \leqslant \frac{4}{3}a_{n+1}^4 + 6a_{n+1}^2M_n^2 + M_n^4 \text{ a.s.}$$

Taking expectation on both sides, we find that

$$\mathbb{E}[M_{n+1}^4] \leqslant \frac{4}{3}a_{n+1}^4 + 6a_{n+1}^2 \mathbb{E}[M_n^2] + \mathbb{E}[M_n^4]$$

leading, via (B.11), to

$$\mathbb{E}[M_n^4] \leqslant \frac{4}{3} \sum_{k=1}^n a_k^4 + 6 \sum_{k=1}^n a_k^2 \mathbb{E}[M_{k-1}^2] \leqslant 6 \left(1 + \sum_{k=1}^n a_k^2\right) \sum_{k=1}^n a_k^2.$$

Finally, we can deduce from (A.10) that

$$\sup_{n\geqslant 1} \mathbb{E}[M_n^4] < \infty,\tag{B.15}$$

which completes the proof of theorem 3.7.

The proof of theorem 3.8 relies on the following well-known lemma on sums of ratio of gamma functions.

Lemma B.1. For any non-negative real numbers a and b such that $b \neq a + 1$ and for all $n \geqslant 1$, we have

$$\sum_{k=1}^{n} \frac{\Gamma(k+a)}{\Gamma(k+b)} = \frac{\Gamma(n+a+1)}{(b-a-1)\Gamma(n+b)} \left(\frac{\Gamma(n+b)\Gamma(a+1)}{\Gamma(n+a+1)\Gamma(b)} - 1 \right). \tag{B.16}$$

Proof of theorem 3.8. Denote $\alpha = 2p - 1$, $\beta = 2q - 1$,

$$L_n = rac{M_n}{\Gamma(2p)}$$
 and $L = rac{M}{\Gamma(2p)}$.

It follows from convergence (3.12) that for any integer d = 1, ..., 4

$$\lim_{n \to \infty} \mathbb{E}[L_n^d] = \mathbb{E}[L^d]. \tag{B.17}$$

First of all, we already saw that for all $n \ge 1$, $\mathbb{E}[S_{n+1}|\mathcal{F}_n] = \gamma_n S_n$ a.s. Consequently,

$$\mathbb{E}[S_{n+1}] = \gamma_n \mathbb{E}[S_n] = \left(\frac{n+\alpha}{n}\right) \mathbb{E}[S_n],$$

which leads to

$$\mathbb{E}[S_n] = \prod_{k=1}^{n-1} \left(\frac{k+\alpha}{k}\right) \mathbb{E}[S_1] = \frac{\beta \Gamma(n+\alpha)}{\Gamma(n)\Gamma(\alpha+1)} = \frac{\beta}{a_n}.$$
 (B.18)

Hence, we immediately get from (B.18) that

$$\mathbb{E}[L_n] = \mathbb{E}[L] = \frac{\beta}{\Gamma(\alpha+1)} = \frac{2q-1}{\Gamma(2p)}.$$

Next, taking expectation on both sides of (A.3), we obtain that for all $n \ge 1$,

$$\mathbb{E}[S_{n+1}^2] = 1 + (2\gamma_n - 1)\mathbb{E}[S_n^2] = 1 + \left(\frac{n + 2\alpha}{n}\right)\mathbb{E}[S_n^2]$$

which implies that

$$\mathbb{E}[S_n^2] = \frac{\Gamma(n+2\alpha)}{\Gamma(n)\Gamma(2\alpha+1)} \left(1 + \sum_{k=1}^{n-1} \frac{\Gamma(k+1)\Gamma(2\alpha+1)}{\Gamma(k+2\alpha+1)} \right),$$
$$= \frac{\Gamma(n+2\alpha)}{\Gamma(n)} \sum_{k=1}^{n} \frac{\Gamma(k)}{\Gamma(k+2\alpha)}.$$

Therefore, we deduce from identity (B.16) with a = 0 and $b = 2\alpha$ that

$$\mathbb{E}[S_n^2] = \frac{n}{2\alpha - 1} \left(\frac{\Gamma(n + 2\alpha)}{\Gamma(n+1)\Gamma(2\alpha)} - 1 \right). \tag{B.19}$$

Hence, we obtain from (B.19) that

$$\mathbb{E}[L_n^2] = \frac{a_n^2 \mathbb{E}[S_n^2]}{(\Gamma(2p))^2} = \frac{n}{2\alpha - 1} \left(\frac{\Gamma(n)}{\Gamma(n + \alpha)}\right)^2 \left(\frac{\Gamma(n + 2\alpha)}{\Gamma(n + 1)\Gamma(2\alpha)} - 1\right)$$

which ensures, via (B.17) with d = 2, that

$$\lim_{n \to \infty} \mathbb{E}[L_n^2] = \mathbb{E}[L^2] = \frac{1}{(2\alpha - 1)\Gamma(2\alpha)} = \frac{1}{(4p - 3)\Gamma(2(2p - 1))}.$$

Furthermore, as $S_{n+1}^3 = (S_n + X_{n+1})^3$, we obtain that

$$\mathbb{E}[S_{n+1}^3 | \mathcal{F}_n] = (3\gamma_n - 2)S_n^3 + (\gamma_n + 2)S_n$$
a.s.

Consequently,

$$\mathbb{E}[S_{n+1}^3] = \left(\frac{3n+\alpha}{n}\right)\mathbb{E}[S_n] + \left(\frac{n+3\alpha}{n}\right)\mathbb{E}[S_n^3].$$

Hence, it follows from tedious but straighforward calculations that

$$\mathbb{E}[S_n^3] = \frac{\beta \Gamma(n+3\alpha)}{\Gamma(n)\Gamma(3\alpha+1)} \left(1 + \frac{3\Gamma(3\alpha+1)}{\Gamma(\alpha+1)} \xi_n \right)$$
 (B.20)

where

$$\xi_n = \sum_{k=1}^{n-1} \left(k + \frac{\alpha}{3}\right) \frac{\Gamma(k+\alpha)}{\Gamma(k+3\alpha+1)} = \sum_{k=1}^{n-1} \frac{\Gamma(k+\alpha+1)}{\Gamma(k+3\alpha+1)} - \frac{2\alpha}{3} \sum_{k=1}^{n-1} \frac{\Gamma(k+\alpha)}{\Gamma(k+3\alpha+1)}.$$

However, we infer from (B.16) with $a = \alpha$ or $a = \alpha + 1$ and $b = 3\alpha + 1$ that

$$\xi_n = \frac{1}{2\alpha - 1} \left(\frac{\Gamma(\alpha + 2)}{\Gamma(3\alpha + 1)} - \frac{\Gamma(n + \alpha + 1)}{\Gamma(n + 3\alpha)} \right) - \frac{1}{3} \left(\frac{\Gamma(\alpha + 1)}{\Gamma(3\alpha + 1)} - \frac{\Gamma(n + \alpha)}{\Gamma(n + 3\alpha)} \right).$$

Therefore, we obtain from (B.20) that

$$\mathbb{E}[S_n^3] = \frac{\beta \Gamma(n+3\alpha)}{\Gamma(n)\Gamma(3\alpha+1)} \left(\frac{3(\alpha+1)}{2\alpha-1} - \frac{\Gamma(3\alpha+1)\Gamma(n+\alpha)}{\Gamma(\alpha+1)\Gamma(n+3\alpha)} \frac{3n+\alpha+1}{2\alpha-1} \right),$$

$$= \frac{\beta}{(2\alpha-1)\Gamma(n)} \left(\frac{3(\alpha+1)\Gamma(n+3\alpha)}{\Gamma(3\alpha+1)} - \frac{\Gamma(n+\alpha)}{\Gamma(\alpha+1)} (3n+\alpha+1) \right). \tag{B.21}$$

Consequently, we obtain from (B.21) that

$$\mathbb{E}[L_n^3] = \frac{\beta(\Gamma(n))^2}{(2\alpha - 1)(\Gamma(n + \alpha))^3} \left(\frac{3(\alpha + 1)\Gamma(n + 3\alpha)}{\Gamma(3\alpha + 1)} - \frac{\Gamma(n + \alpha)}{\Gamma(\alpha + 1)} (3n + \alpha + 1) \right)$$

which leads, via (B.17) with d = 3, to

$$\lim_{n \to \infty} \mathbb{E}[L_n^3] = \mathbb{E}[L^3] = \frac{\beta(\alpha + 1)}{\alpha(2\alpha - 1)\Gamma(3\alpha)} = \frac{2p(2q - 1)}{(2p - 1)(4p - 3)\Gamma(3(2p - 1))}.$$

By the same token, since $S_{n+1}^4 = (S_n + X_{n+1})^4$, we obtain that

$$\mathbb{E}[S_{n+1}^4] = 1 + 2\left(\frac{3n+2\alpha}{n}\right)\mathbb{E}[S_n^2] + \left(\frac{n+4\alpha}{n}\right)\mathbb{E}[S_n^4],$$

which implies that

$$\mathbb{E}[S_n^4] = \frac{\Gamma(n+4\alpha)}{\Gamma(n)\Gamma(4\alpha+1)} \left(1 + \frac{\Gamma(4\alpha+1)}{(2\alpha-1)\Gamma(2\alpha)} \left(P_n + 6Q_n - 6R_n \right) \right)$$
 (B.22)

where

$$P_{n} = (2\alpha - 1)\Gamma(2\alpha) \sum_{k=1}^{n-1} \frac{\Gamma(k+1)}{\Gamma(k+4\alpha+1)},$$

$$Q_{n} = \sum_{k=1}^{n-1} \left(k + \frac{2\alpha}{3}\right) \frac{\Gamma(k+2\alpha)}{\Gamma(k+4\alpha+1)},$$

$$R_{n} = \Gamma(2\alpha) \sum_{k=1}^{n-1} \left(k + \frac{2\alpha}{3}\right) \frac{\Gamma(k+1)}{\Gamma(k+4\alpha+1)}.$$

We make use once again of identity (B.16) with appropriate values of a and b, to find that

$$\mathbb{E}[S_n^4] = \frac{\Gamma(n+4\alpha)}{\Gamma(n)\Gamma(4\alpha+1)} \left(\frac{24\alpha(2\alpha(\alpha+1)-1)}{(2\alpha-1)^2(4\alpha-1)} - \frac{\Gamma(4\alpha+1)}{(2\alpha-1)^2\Gamma(n+4\alpha)} \zeta_n \right)$$

where

$$\zeta_n = \frac{2(3n+2(\alpha+1))}{\Gamma(2\alpha)}\Gamma(n+2\alpha) - \frac{3n(4\alpha-1)+2(2\alpha^2+1)}{(4\alpha-1)}\Gamma(n+1).$$

Finally, we deduce from (B.17) with d = 4 that

$$\begin{split} \lim_{n\to\infty} \mathbb{E}[L_n^4] &= \mathbb{E}[L^4] = \frac{24\alpha(2\alpha(\alpha+1)-1)}{(2\alpha-1)^2(4\alpha-1)\Gamma(4\alpha+1)} = \frac{6(2\alpha(\alpha+1)-1)}{(2\alpha-1)^2(4\alpha-1)\Gamma(4\alpha)}, \\ &= \frac{6(8p^2-4p-1)}{(8p-5)(4p-3)^2\Gamma(4(2p-1))}, \end{split}$$

which completes the proof of theorem 3.8.

Appendix C. Proofs of the asymptotic normality results

C.1. The diffusive regime

Proof of theorem 3.3. We shall make use of the central limit theorem for martingales given, e.g., by corollary 2.1.10 of [9], to establish the asymptotic normality (3.5). It follows from (3.1), (A.7) and (A.8) that

$$\lim_{n \to \infty} \frac{\langle M \rangle_n}{v_n} = 1 \text{ a.s.}$$
 (C.1)

Hereafter, it only remains to prove that (M_n) satisfies Lindeberg's condition, that is, for all $\varepsilon > 0$,

$$\frac{1}{\nu_n} \sum_{k=1}^n \mathbb{E}\left[|\Delta M_k|^2 \mathbf{I}_{|\Delta M_k| \geqslant \varepsilon \sqrt{\nu_n}} \middle| \mathcal{F}_{k-1} \right] \xrightarrow{\mathcal{P}} 0. \tag{C.2}$$

We obtain from (A.6) that for all $\varepsilon > 0$,

$$\frac{1}{v_n} \sum_{k=1}^n \mathbb{E} \left[|\Delta M_k|^2 \mathbf{I}_{|\Delta M_k| \geqslant \varepsilon \sqrt{v_n}} | \mathcal{F}_{k-1} \right] \leqslant \frac{1}{\varepsilon^2 v_n^2} \sum_{k=1}^n \mathbb{E} \left[|\Delta M_k|^4 | \mathcal{F}_{k-1} \right],$$

$$\leqslant \sup_{1 \leqslant k \leqslant n} \mathbb{E} \left[\varepsilon_k^4 | \mathcal{F}_{k-1} \right] \frac{1}{\varepsilon^2 v_n^2} \sum_{k=1}^n a_k^4,$$

$$\leqslant \frac{4}{3 \varepsilon^2 v_n^2} \sum_{k=1}^n a_k^4.$$

However, we deduce from (B.5), together with Kronecker's lemma, that

$$\lim_{n\to\infty}\frac{1}{v_n^2}\sum_{k=1}^n a_k^4=0,$$

which ensures that Lindeberg's condition is satisfied. Therefore, we can conclude from the central limit theorem for martingales that

$$\frac{1}{\sqrt{v_n}} M_n \xrightarrow{\mathcal{L}} \mathcal{N}(0, 1). \tag{C.3}$$

As $M_n = a_n S_n$ and $\sqrt{n} a_n$ is equivalent to $\sqrt{v_n(3-4p)}$, we find from (C.3) that

$$\frac{S_n}{\sqrt{n}} \xrightarrow{\mathcal{L}} \mathcal{N}\left(0, \frac{1}{3-4p}\right),$$

which is exactly what we wanted to prove.

C.2. The critical regime

Proof of theorem 3.6. We shall now proceed to the proof of the asymptotic normality (3.10) making use, once again, of the central limit theorem for martingales given e.g. by corollary 2.1.10 of [9]. We have from (3.6), (A.7) and (A.9) that

$$\lim_{n \to \infty} \frac{\langle M \rangle_n}{v_n} = 1 \text{ a.s.}$$
 (C.4)

Moreover, it follows from (A.9) and (B.8) that $a_n^2 v_n^{-1}$ is equivalent to $(n \log n)^{-1}$. However, it is well known that

$$\sum_{n=1}^{+\infty} \frac{1}{n^2 (\log n)^2} < \infty.$$

Consequently, we clearly have

$$\sum_{n=1}^{+\infty} \frac{a_n^4}{v_n^2} < +\infty. \tag{C.5}$$

We already saw in the proof of theorem 3.3 that (C.5) is enough to check that (M_n) satisfies Lindeberg's condition. Consequently, we can deduce from the central limit theorem for martingales that

$$\frac{1}{\sqrt{\nu_n}} M_n \xrightarrow{\mathcal{L}} \mathcal{N}(0,1). \tag{C.6}$$

Hence, as $M_n = a_n S_n$ and $\sqrt{n \log n} a_n$ is equivalent to $\sqrt{v_n}$, we find from (C.6) that

$$\frac{S_n}{\sqrt{n\log n}} \xrightarrow{\mathcal{L}} \mathcal{N}(0,1),$$

which achieves the proof of theorem 3.6.

ORCID iDs

Bernard Bercu https://orcid.org/0000-0002-8118-5232

References

- Baur E and Bertoin J 2016 Elephant random walks and their connection to Pólya-type urns *Phys. Rev.* E 94 052134
- [2] Boyer D and Romo-Cruz J C R 2014 Solvable random-walk model with memory and its relations with Markovian models of anomalous diffusion *Phys. Rev.* E **90** 042136
- [3] Bercu B 2004 On the convergence of moments in the almost sure central limit theorem for martingales with statistical applications Stoch. Process. Appl. 111 157–73
- [4] Chauvin B, Pouyanne N and Sahnoun R 2011 Limit distributions for large Pólya urns Ann. Appl. Probab. 21 1–32
- [5] Coletti C F, Gava R and Schütz G M 2017 Central limit theorem and related results for the elephant random walk J. Math. Phys. 58 053303
- [6] Cressoni J C, Viswanathan G M and Da Silva M A A 2013 Exact solution of an anisotropic 2D random walk model with strong memory correlations J. Phys. A: Math. Theor. 46 505002
- [7] Cressoni J C, Da Silva M A A and Viswanathan G M 2007 Amnestically induced persistence in random walks Phys. Rev. Lett. 98 070603
- [8] Da Silva M A A, Cressoni J C, Schütz G M, Viswanathan G M and Trimper S 2013 Non-Gaussian propagator for elephant random walks *Phys. Rev.* E 88 022115

- [9] Duflo M 1997 Random Iterative Models (Applications of Mathematics vol 34) (Berlin: Springer)
- [10] Janson S 2004 Functional limit theorems for multitype branching processes and generalized Pólya urns Stoch. Process. Appl. 110 177–245
- [11] Kumar N, Harbola U and Lindenberg K 2010 Memory-induced anomalous dynamics: emergence of diffusion, subdiffusion, and superdiffusion from a single random walk model *Phys. Rev.* E 82 021101
- [12] Hall P and Heyde C C 1980 Martingale Limit Theory and its Application (New York: Academic)
- [13] Harbola U, Kumar N and Lindenberg K 2014 Memory-induced anomalous dynamics in a minimal random walk model *Phys. Rev.* E 90 022136
- [14] Kürsten R 2016 Random recursive trees and the elephant random walk *Phys. Rev.* E 93 032111
- [15] Lyu J, Xin J and Yu Y 2017 Residual diffusivity in elephant random walk models with stops (arXiv:1705.02711)
- [16] Paraan F N C and Esguerra J P 2006 Exact moments in a continuous time random walk with complete memory of its history Phys. Rev. E 74 032101
- [17] Pemantle R 2007 A survey of random processes with reinforcement *Prob. Surv.* 4 1–79
- [18] Schütz G M and Trimper S 2004 Elephants can always remember: exact long-range memory effects in a non-Markovian random walk *Phys. Rev.* E 70 045101
- [19] Stout W F 1970 A martingale analogue of Kolmogorovs law of the iterated logarithm Z. Wahrscheinlichkeitstheor. 15 279–90
- [20] Stout W F 1974 Almost Sure Convergence. Probability, Mathematical Statistics vol 24 (New York: Academic)