# SHARP LARGE DEVIATIONS FOR THE ORNSTEIN-UHLENBECK PROCESS* 

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#### Abstract

We establish sharp large deviation principles for well-known random variables associated with the Ornstein-Uhlenbeck process, such as the energy, the maximum likelihood estimator of the drift parameter, and the log-likelihood ratio.


Key words. large deviations, Ornstein-Uhlenbeck process, likelihood estimation

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1. Introduction. Consider the Ornstein-Uhlenbeck process

$$
\begin{equation*}
d X_{t}=\theta X_{t} d t+d W_{t} \tag{1.1}
\end{equation*}
$$

where $W$ is a standard Brownian motion and the parameter $\theta$ is strictly negative. For the sake of simplicity, we choose the initial state $X_{0}=0$. In this paper, we investigate the sharp large deviation properties for well-known random variables associated with (1.1), such as the energy

$$
\begin{equation*}
S_{T}=\int_{0}^{T} X_{t}^{2} d t \tag{1.2}
\end{equation*}
$$

the maximum likelihood estimator of $\theta$

$$
\begin{equation*}
\hat{\theta}_{T}=\frac{\int_{0}^{T} X_{t} d X_{t}}{\int_{0}^{T} X_{t}^{2} d t}=\frac{X_{T}^{2}-T}{2 \int_{0}^{T} X_{t}^{2} d t} \tag{1.3}
\end{equation*}
$$

and the log-likelihood ratio

$$
\begin{equation*}
V_{T}=\left(\theta_{0}-\theta_{1}\right) \int_{0}^{T} X_{t} d X_{t}-\frac{1}{2}\left(\theta_{0}^{2}-\theta_{1}^{2}\right) \int_{0}^{T} X_{t}^{2} d t \tag{1.4}
\end{equation*}
$$

with $\theta_{0}$ and $\theta_{1}$ strictly negative. It was already proven that a.s., as $T$ goes to infinity,

$$
\begin{equation*}
\frac{S_{T}}{T} \rightarrow-\frac{1}{2 \theta}, \quad \hat{\theta}_{T} \rightarrow \theta, \quad \frac{V_{T}}{T} \rightarrow-\frac{\left(\theta_{0}-\theta_{1}\right)^{2}}{4 \theta_{0}} \tag{1.5}
\end{equation*}
$$

Fluctuations are also known [1]. More recently, Bryc and Dembo [8] and FlorensLandais and Pham [14] have established large deviation principles for $S_{T}$ and $\hat{\theta}_{T}$. Strictly speaking, $\hat{\theta}_{T}$ is not the maximum likelihood estimator of $\theta$ since $\hat{\theta}_{T}$ may take nonnegative values, whereas the parameter $\theta$ is assumed to be strictly negative.

[^0]Nevertheless, we use this terminology throughout the paper. We present here sharp large deviation principles for the random variables $S_{T}, \hat{\theta}_{T}$, and $V_{T}$. As usual, we shall say that a family of real random variables $\left(Z_{T}\right)$ satisfies a large deviation principle (LDP), with rate function $I$, if $I$ is a lower semicontinuous function from $\mathbf{R}$ to $[0,+\infty]$ such that, for any closed set $F \subset \mathbf{R}, \lim \sup _{T \rightarrow \infty} T^{-1} \log \mathbf{P}\left\{Z_{T} \in F\right\} \leqq-\inf _{x \in F} I(x)$, while for any open set $G \subset \mathbf{R}$,

$$
-\inf _{x \in G} I(x) \leqq \liminf _{T \rightarrow \infty} \frac{1}{T} \log \mathbf{P}\left\{Z_{T} \in G\right\}
$$

Moreover, $I$ is a good rate function if its level sets are compact subsets of $\mathbf{R}$. When $I$ has a unique minimum $m$, which will always be the case here, an LDP for $\left(Z_{T}\right)$ gives the asymptotic behavior of $\mathbf{P}\left\{Z_{T} \geqq c\right\}$ or $\mathbf{P}\left\{Z_{T} \leqq c\right\}$ in a logarithmic scale whenever $c>m$ or $c<m$, respectively. We shall say that a sequence of real random variables $\left(Z_{T}\right)$ satisfies a sharp large deviation principle (SLDP) if, for any real number $c$, it is possible to give asymptotic expansions of $e^{T I(c)} \mathbf{P}\left\{Z_{T} \geqq c\right\}$ or $e^{T I(c)} \mathbf{P}\left\{Z_{T} \leqq c\right\}$ in powers of $T^{-1}$. For the sake of simplicity, we present only the first one.

In order to prove an LDP for (1.2), (1.3), or (1.4), the main tool is the normalized cumulant generating function (c.g.f.) of the pair $\left(\int_{0}^{T} X_{t} d X_{t}, \int_{0}^{T} X_{t}^{2} d t\right)$

$$
\begin{equation*}
\mathcal{L}_{T}(a, b)=\frac{1}{T} \log \mathbf{E}\left[\exp \left(\mathcal{Z}_{T}(a, b)\right)\right] \tag{1.6}
\end{equation*}
$$

where, for any $(a, b) \in \mathbf{R}^{2}$,

$$
\begin{equation*}
\mathcal{Z}_{T}(a, b)=a \int_{0}^{T} X_{t} d X_{t}+b \int_{0}^{T} X_{t}^{2} d t \tag{1.7}
\end{equation*}
$$

It is possible to establish SLDP for all linear combinations of $\int_{0}^{T} X_{t} d X_{t}$ and $\int_{0}^{T} X_{t}^{2} d t$. However, in order to improve the presentation of the paper, we prefer to focus our attention on the random variables (1.2), (1.3), and (1.4).

In discrete time, the analogue of (1.1) is the first order autoregressive process. It was studied, as a particular case, in [4] for LDP and in [5] for SLDP. The proofs used Toeplitz matrices and their asymptotic spectral properties given by the first Szegö theorem for LDP and the strong Szegö theorem for SLDP.

The covariance of the stationary Ornstein-Uhlenbeck process is a Wiener-Hopf operator and, for some "quadratic forms," we can use a similar scheme [23]. Here, we take advantage of the explicit expression of the c.g.f.

An important point for proving an LDP is the determination of the domain $\Delta$ of the limit $\mathcal{L}$ of the c.g.f. We shall say that $\mathcal{L}$ is steep if its derivative has an infinite limit at the boundary of $\Delta$ (see, e.g., [9, p. 44]). It is a sufficient condition to apply the Gärtner-Ellis theorem. The main difficulty arises when $\mathcal{L}$ is not steep.

Before entering into details of the different cases, let us summarize the general scheme. We have to study probabilities of the form $\mathbf{P}\left\{\mathcal{Z}_{T}(a, b)>z T\right\}$. We follow a similar approach to the one recently given in [5] for discrete time quadratic forms of Gaussian stationary processes. It was inspired by the original work of Bahadur and Rao [2] for the sample mean of a sequence of independent and identically distributed (i.i.d.) random variables. We perform an exponential change of probability of parameter $\varphi$ to be chosen to "track" $z$,

$$
\begin{equation*}
\frac{d \mathbf{Q}_{T}}{d \mathbf{P}}=\exp \left(\varphi \mathcal{Z}_{T}(a, b)-T \mathcal{L}_{T}(a \varphi, b \varphi)\right) \tag{1.8}
\end{equation*}
$$

Then, we have the decomposition

$$
\mathbf{P}\left\{\mathcal{Z}_{T}(a, b)>z T\right\}=A_{T} B_{T}
$$

with

$$
\begin{align*}
A_{T} & =\exp \left[T\left(\mathcal{L}_{T}(a \varphi, b \varphi)-z \varphi\right)\right]  \tag{1.9}\\
B_{T} & =\mathbf{E}_{T}\left(\exp \left[-\varphi\left(\mathcal{Z}_{T}(a, b)-z T\right)\right] \mathbb{I}_{\mathcal{Z}_{T}(a, b) \geqq z T}\right) \tag{1.10}
\end{align*}
$$

where $\mathbf{E}_{T}$ is the expectation under $\mathbf{Q}_{T}$. On the one hand, an asymptotic expansion for $A_{T}$ is given in section 2. On the other hand, we expand $B_{T}$ via an evaluation of the characteristic function $\Phi_{T}$ of $\left(\mathcal{Z}_{T}(a, b)-z T\right)$ (properly normalized) under the new probability $\mathbf{Q}_{T}$. In general, $\Phi_{T}$ has a pointwise limit. This is, however, not enough to study $B_{T}$. Actually, we obtain a complete expansion of $\Phi_{T}$ similar to the one established in the i.i.d. case by Cramer [7] and Esseen [12] and we integrate it using Parseval's theorem. This approach was developed in other contexts with different asymptotics by Ben Arous [3], Bolthausen [6], Dembo, Mayer-Wolf, and Zeitouni [11], Ibragimov [16], Li [19], and Zolotarev [24] (see also the references therein).

In the steep case, the limit distribution is $N(0,1)$ and the convergence of $\Phi_{T}$ is dominated. Thus, we get an SLDP similar to the one obtained in discrete time [5]. In particular, for $c>-1 /(2 \theta)$, we prove that $\sqrt{T} e^{T I(c)} \mathbf{P}\left\{S_{T} \geqq c T\right\}$ has a limit given explicitly.

In the nonsteep case, occurring only for (1.3), we find new regimes. For $c \geqq \theta / 3$, the rate function $I$ is linear. In the change of probability (1.8), we use a time varying parameter $\varphi_{T}$. This strategy was developed in [10] and [8]. Under this new probability, for $c>\theta / 3$, the limit distribution is a centered $\chi^{2}$ (see also [14]). To establish an SLDP, this gives rise to a new problem since uniform integrability is not guaranteed. Eventually, the SLDP has the same form as in the steep case but with different constants. For $c=\theta / 3, \varphi_{T}$ tends to the boundary of $\Delta$ at a different rate and the limit distribution is the convolution of an $N(0,1)$ and a centered $\chi^{2}$. This yields an SLDP with rate $\sqrt{T}$ instead of $T$. In particular, $T^{1 / 4} e^{T I(c)} \mathbf{P}\left\{\hat{\theta}_{T} \geqq c\right\}$ has a limit given explicitly.

The paper is organized as follows. In section 2, we give a sharp description of $\mathcal{L}_{T}(a, b)$ and $\mathcal{Z}_{T}(a, b)$. Then, we present the SLDP results: section 3 is devoted to the energy, section 4 to the maximum likelihood estimator, and section 5 to the log-likelihood ratio. Proofs are collected in sections 6 and 7.
2. Main tools. Before presenting specific results for each functional, we give two lemmas which are the core of all our proofs and which allow unified notation.

The convergence of (1.6) was studied in [14]. Lemma 2.1 below gives a new presentation enlightening the role of the limit $\mathcal{L}$ and of the first order term $\mathcal{H}$ for both LDP and SLDP. This decomposition is completely analogous to the one given in discrete time [5].

It is well known that c.g.f.'s of quadratic functionals of Gaussian processes are connected to empirical distributions of eigenvalues of some operator (see, e.g., [15, Chap. 11]).

In our case, Lemma 2.2 gives a decomposition of $\mathcal{Z}_{T}$ using eigenvalues (see section 7.1 for the operator) and a convergence of their empirical distribution. This last part is similar to the first Szegö theorem and uses the convergence of $\mathcal{L}_{T}$ to $\mathcal{L}$ for its proof.

The limit is proportional to the image of the Lebesgue measure by the spectral density of the stationary Ornstein-Uhlenbeck process of parameter $\theta<0$

$$
g(x)=\frac{1}{\theta^{2}+x^{2}}
$$

associated with the covariance $r(t)=-\exp (\theta|t|) /(2 \theta)$.
Lemma 2.1. Set $\Delta=\left\{(a, b) \in \mathbf{R}^{2} \mid \theta^{2}-2 b>0\right.$ and $\left.\theta+a<\sqrt{\theta^{2}-2 b}\right\}$ and let $\rho(b)=\sqrt{\theta^{2}-2 b}$.
(i) For all $(a, b) \in \Delta$, we have

$$
\begin{equation*}
\mathcal{L}_{T}(a, b)=\mathcal{L}(a, b)+\frac{\mathcal{H}(a, b)}{T}+\frac{1}{T} \mathcal{R}_{T}(a, b) \tag{2.1}
\end{equation*}
$$

with

$$
\begin{align*}
& \mathcal{L}(a, b)=-\frac{1}{2}(a+\theta+\rho(b)) \\
& \mathcal{H}(a, b)=-\frac{1}{2} \log \left(\frac{1}{2}\left(1-(a+\theta) \rho^{-1}(b)\right)\right)  \tag{2.2}\\
& \mathcal{R}_{T}(a, b)=-\frac{1}{2} \log \left(1+\frac{1+(a+\theta) \rho^{-1}(b)}{1-(a+\theta) \rho^{-1}(b)} e^{-2 T \rho(b)}\right) . \tag{2.3}
\end{align*}
$$

(ii) Moreover, the remainder $\mathcal{R}_{T}(a, b)$ goes exponentially fast to zero:

$$
\begin{equation*}
\mathcal{R}_{T}(a, b)=\mathcal{O}\left(e^{-2 T \rho(b)}\right) \tag{2.4}
\end{equation*}
$$

Denote by $\mathcal{F}$ the class of all continuous functions $f$ on $\mathbf{R}$ such that $f(x)=x h(x)$ with $h$ continuous.

Lemma 2.2. (i) There exists a sequence of real numbers $\left(\lambda_{j}^{T}\right)$ such that $\left(\lambda_{j}^{T}\right) \in$ $l^{1}(\mathbf{N})$,

$$
\begin{equation*}
\mathcal{Z}_{T}(a, b)=-\frac{a T}{2}+\sum_{j=1}^{\infty} \lambda_{j}^{T} \varepsilon_{j}^{2} \tag{2.5}
\end{equation*}
$$

where $\left(\varepsilon_{j}\right)$ are independent standard $N(0,1)$ random variables.
(ii) There exists a fixed compact $[A, B]$ such that $\lambda_{j}^{T} \in[A, B]$ for every $j$ and $T$.
(iii) If $b \neq 0$, then the empirical measure

$$
\begin{equation*}
\nu_{T}=\frac{1}{T} \sum_{j=1}^{\infty} \delta_{\lambda_{j}^{T}} \tag{2.6}
\end{equation*}
$$

converges for the duality with functions in $\mathcal{F}$ to $\nu$ defined, for any continuous function $h$ with compact support, by

$$
\begin{equation*}
\langle\nu, h\rangle=\frac{1}{2 \pi} \int_{\mathbf{R}} h(b g(x)) d x \tag{2.7}
\end{equation*}
$$

Remarks. (1) The sequence $\left(\lambda_{j}^{T}\right)$ is finite if $b=0$ and infinite otherwise.
(2) The case $a=0$ is already known. Fix $b=1$ to simplify. The KarhunenLoève expansion of the process $\left(X_{t}\right)$ directly gives (i). Bryc and Dembo [8] proved (ii) and (iii) (see also [15, Chaps. 8 and 11]). In this case, $A=0$ and $B=\theta^{-2}$.
(3) If $a$ and $b$ are both nonnegative (respectively, nonpositive), all the $\left(\lambda_{j}^{T}\right)$ are nonnegative and we may take $A=0$ (respectively, $B=0$ ).
(4) In the particular case of the likelihood ratio (1.4), decomposition (2.5) is well known (see, e.g., [18] and [22]).

Warning. The functions $I, L, L_{T}$ given in the following sections are different. They are defined at the beginning of each section. In order to avoid heaviness in the notation, we choose to keep the same symbol for quantities of the same nature.
3. Energy. The LDP for $\left(S_{T}\right)$ was proved by Bryc and Dembo [8] for general centered stationary Gaussian processes. In the particular case of the Ornstein-Uhlenbeck process (1.1), they have the following result.

Lemma 3.1. $\left(T^{-1} S_{T}\right)$ satisfies an $L D P$ with good rate function

$$
I(c)= \begin{cases}\frac{(2 \theta c+1)^{2}}{8 c} & \text { if } c>0  \tag{3.1}\\ +\infty & \text { otherwise }\end{cases}
$$

The rate function $I$ has a unique minimum in $c=-1 /(2 \theta)$ which is the a.s. limit given in (1.5). We are going to improve this result by an SLDP for $\left(S_{T}\right)$ similar to the well-known Bahadur-Rao theorem [2]. Set, for all $a<\theta^{2} / 2$,

$$
\begin{equation*}
L(a)=\mathcal{L}(0, a) \quad \text { and } \quad H(a)=\mathcal{H}(0, a) \tag{3.2}
\end{equation*}
$$

One can remark that the rate function $I$ given by (3.1) is the Fenchel-Legendre dual of the function $L$.

ThEOREM 3.1. $\left(T^{-1} S_{T}\right)$ satisfies an SLDP associated with $L$ and $H$. More precisely, for all $c>-1 /(2 \theta)$, there exists a sequence $\left(b_{c, k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
\mathbf{P}\left\{S_{T} \geqq c T\right\}=\frac{e^{-T I(c)+H\left(a_{c}\right)}}{\sigma_{c} a_{c} \sqrt{2 \pi T}}\left[1+\sum_{k=1}^{p} \frac{b_{c, k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right] \tag{3.3}
\end{equation*}
$$

with

$$
\begin{equation*}
a_{c}=\frac{4 \theta^{2} c^{2}-1}{8 c^{2}}, \quad H\left(a_{c}\right)=-\frac{1}{2} \log \left(\frac{1}{2}(1-2 \theta c)\right) \tag{3.4}
\end{equation*}
$$

and $\sigma_{c}^{2}=L^{\prime \prime}\left(a_{c}\right)=4 c^{3}$. The coefficients $b_{c, 1}, b_{c, 2}, \ldots, b_{c, p}$ may be explicitly given as functions of the derivatives of $L$ and $H$ at point $a_{c}$. For example, the first coefficient $b_{c, 1}$ is given by

$$
\begin{equation*}
b_{c, 1}=\frac{1}{\sigma_{c}^{2}}\left(-\frac{h_{2}}{2}-\frac{h_{1}^{2}}{2}+\frac{l_{4}}{8 \sigma_{c}^{2}}+\frac{l_{3} h_{1}}{2 \sigma_{c}^{2}}-\frac{5 l_{3}^{2}}{24 \sigma_{c}^{4}}+\frac{h_{1}}{a_{c}}-\frac{l_{3}}{2 a_{c} \sigma_{c}^{2}}-\frac{1}{a_{c}^{2}}\right) \tag{3.5}
\end{equation*}
$$

with $l_{k}=L^{(k)}\left(a_{c}\right)$ and $h_{k}=H^{(k)}\left(a_{c}\right)$. More precisely,

$$
\begin{equation*}
b_{c, 1}=\frac{-2 c\left(12 \theta^{4} c^{4}+12 \theta^{3} c^{3}+35 \theta^{2} c^{2}+4 \theta c+2\right)}{\left(4 \theta^{2} c^{2}-1\right)^{2}} \tag{3.6}
\end{equation*}
$$

4. Maximum likelihood estimator. Florens-Landais and Pham [14] proved the following LDP for $\left(\hat{\theta}_{T}\right)$.

Lemma 4.1. $\left(\hat{\theta}_{T}\right)$ satisfies an LDP with good rate function

$$
I(c)= \begin{cases}-\frac{(c-\theta)^{2}}{4 c} & \text { if } c<\frac{\theta}{3}  \tag{4.1}\\ 2 c-\theta & \text { otherwise }\end{cases}
$$

The rate function $I$ has a unique minimum in $c=\theta$ which is the a.s. limit given in (1.5). The large deviation properties of $\left(\hat{\theta}_{T}\right)$ are related to the ones of

$$
\begin{equation*}
Z_{T}(c)=\int_{0}^{T} X_{t} d X_{t}-c \int_{0}^{T} X_{t}^{2} d t \tag{4.2}
\end{equation*}
$$

with $c \in \mathbf{R}$ since $\mathbf{P}\left\{\hat{\theta}_{T} \geqq c\right\}=\mathbf{P}\left\{Z_{T}(c) \geqq 0\right\}$. One has to keep in mind that the threshold $c$ for $\hat{\theta}_{T}$ appears as a parameter for $Z_{T}(c)$. As before, we also improve Lemma 4.1 by an SLDP for $\left(\hat{\theta}_{T}\right)$. Define $\Gamma=\left\{a \in \mathbf{R} \mid \theta^{2}+2 a c>0\right.$ and $\theta+a<$ $\left.\sqrt{\theta^{2}+2 a c}\right\}$ and set, for all $a \in \Gamma$,

$$
\begin{equation*}
L(a)=\mathcal{L}(a,-c a) \quad \text { and } \quad H(a)=\mathcal{H}(a,-c a) \tag{4.3}
\end{equation*}
$$

The rate function $I$ given by (4.1) is $I(c)=-\inf _{a} L(a)$. It is easy to check that

$$
\Gamma= \begin{cases}]-\infty,-\frac{\theta^{2}}{2 c}[ & \text { if } c \leqq \frac{\theta}{2},  \tag{4.4}\\ ]-\infty, a^{c}[ & \text { if } \frac{\theta}{2}<c \leqq 0, \\ ]-\frac{\theta^{2}}{2 c}, a^{c}[ & \text { if } c>0\end{cases}
$$

with $a^{c}=2(c-\theta)$. The main difficulty in comparing this with the previous section is that $L$ is not always steep. Actually, for all $a \in \Gamma$

$$
\begin{equation*}
L(a)=-\frac{1}{2}\left(a+\theta+\sqrt{\theta^{2}+2 a c}\right), \quad L^{\prime}(a)=-\frac{1}{2}\left(1+\frac{c}{\sqrt{\theta^{2}+2 a c}}\right) \tag{4.5}
\end{equation*}
$$

Consequently, for $c \leqq \theta / 2, L$ is steep while this is no longer true for $c>\theta / 2$ since $L^{\prime}\left(a^{c}\right)=-\frac{1}{2}((3 c-\bar{\theta}) /(2 c-\theta))$. Moreover, $L^{\prime}(a)=0$ if and only if $a=a_{c}$ with $a_{c}=\left(c^{2}-\theta^{2}\right) /(2 c)$ and $a_{c} \in \Gamma$ only when $c<\theta / 3$. Therefore, $I(c)=-L\left(a_{c}\right)$ if $c<\theta / 3$ and $I(c)=-L\left(a^{c}\right)$ otherwise.

THEOREM 4.1. $\left(\hat{\theta}_{T}\right)$ satisfies an SLDP associated with $L$ and $H$. More precisely, for all $\theta<c<\theta / 3$, there exists a sequence $\left(b_{c, k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
\mathbf{P}\left\{\hat{\theta}_{T} \geqq c\right\}=\frac{e^{-T I(c)+H\left(a_{c}\right)}}{\sigma_{c} a_{c} \sqrt{2 \pi T}}\left[1+\sum_{k=1}^{p} \frac{b_{c, k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right] \tag{4.6}
\end{equation*}
$$

with

$$
\begin{equation*}
a_{c}=\frac{c^{2}-\theta^{2}}{2 c}, \quad H\left(a_{c}\right)=-\frac{1}{2} \log \frac{(c+\theta)(3 c-\theta)}{4 c^{2}} \tag{4.7}
\end{equation*}
$$

and $\sigma_{c}^{2}=L^{\prime \prime}\left(a_{c}\right)=-(2 c)^{-1}$. The coefficients $b_{c, 1}, b_{c, 2}, \ldots, b_{c, p}$ may be explicitly given as in Theorem 3.1. Furthermore, for $c>\theta / 3$ with $c \neq 0$, there exists a sequence $\left(d_{c, k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
\mathbf{P}\left\{\hat{\theta}_{T} \geqq c\right\}=\frac{e^{-T I(c)+K(c)}}{\sigma^{c} a^{c} \sqrt{2 \pi T}}\left[1+\sum_{k=1}^{p} \frac{d_{c, k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right] \tag{4.8}
\end{equation*}
$$

with $a^{c}=2(c-\theta)$,

$$
\begin{equation*}
\left(\sigma^{c}\right)^{2}=L^{\prime \prime}\left(a^{c}\right)=\frac{c^{2}}{2(2 c-\theta)^{3}}, \quad K(c)=-\frac{1}{2} \log \frac{(c-\theta)(3 c-\theta)}{4 c^{2}} \tag{4.9}
\end{equation*}
$$

The coefficients $d_{c, 1}, d_{c, 2}, \ldots, d_{c, p}$ may be explicitly calculated. For example,

$$
\begin{equation*}
d_{c, 1}=\frac{2 c^{4}-c^{3}(19+3 \theta)+c^{2} \theta(23+\theta)-12 c \theta^{2}+2 \theta^{3}}{4(c-\theta)(2 c-\theta)(3 c-\theta)^{2}} \tag{4.10}
\end{equation*}
$$

Finally, for $c=0, p>0$, and for $T$ large enough,

$$
\begin{equation*}
\mathbf{P}\left\{\hat{\theta}_{T} \geqq 0\right\}=2 \frac{e^{-T I(c)}}{\sqrt{2 \pi T} \sqrt{-2 \theta}}\left[1+\sum_{k=1}^{p} \frac{(2 k)!}{2^{2 k} \theta^{k} T^{k} k!}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right] \tag{4.11}
\end{equation*}
$$

Remark. As $\mathbf{P}\left\{\hat{\theta}_{T} \geqq 0\right\}=\mathbf{P}\left\{X_{T}^{2} \geqq T\right\}$ and $X_{T}$ is Gaussian, (4.11) immediately follows from [13] (see relation (7.1), p. 193). It is easy to check that, for $c \rightarrow 0$, the main part of (4.8) and the first coefficient $d_{c, 1}$ given by (4.10) coincide with the corresponding terms in (4.11).

THEOREM 4.2. For $c=\theta / 3$, there exists a sequence $\left(d_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
\mathbf{P}\left\{\hat{\theta}_{T} \geqq c\right\}=\frac{e^{-T I(c)}}{2 \pi T^{1 / 4}} \frac{\Gamma(1 / 4)}{a_{\theta}^{3 / 4} \sigma_{\theta}}\left[1+\sum_{k=1}^{2 p} \frac{d_{k}}{(\sqrt{T})^{k}}+\mathcal{O}\left(\frac{1}{T^{p} \sqrt{T}}\right)\right] \tag{4.12}
\end{equation*}
$$

with $a_{\theta}=-4 \theta / 3$ and $\sigma_{\theta}^{2}=L^{\prime \prime}\left(a_{\theta}\right)=-3 /(2 \theta)$. As before, the coefficients $d_{1}, d_{2}, \ldots$, $d_{2 p}$ may be explicitly calculated.
5. Log-likelihood ratio. If we wish to test $H_{0}: \theta=\theta_{0}$ versus $H_{1}: \theta=\theta_{1}$ for the Ornstein-Uhlenbeck process (1.1), then the most powerful test is based on the log-likelihood ratio $\left(V_{T}\right)$.

LEMMA 5.1. Under the hypothesis $H_{0},\left(T^{-1} V_{T}\right)$ satisfies an LDP with good rate function

$$
I(c)= \begin{cases}-\frac{1}{8} \frac{\left(\left(\theta_{0}-\theta_{1}\right)^{2}+4 \theta_{0} c\right)^{2}}{\left(\theta_{0}-\theta_{1}+2 c\right)\left(\theta_{0}^{2}-\theta_{1}^{2}\right)} & \text { if } \frac{2 c}{\theta_{0}-\theta_{1}}>-1  \tag{5.1}\\ +\infty & \text { otherwise }\end{cases}
$$

The rate function $I$ has a unique minimum in $c=-\left(\theta_{0}-\theta_{1}\right)^{2} /\left(4 \theta_{0}\right)$ which is the a.s. limit given in (1.5). For all $a \in \mathbf{R}$ such that $\theta_{0}^{2}+\left(\theta_{0}^{2}-\theta_{1}^{2}\right) \times a>0$, set

$$
\begin{align*}
L(a) & =\mathcal{L}\left(a\left(\theta_{0}-\theta_{1}\right),-\frac{1}{2} a\left(\theta_{0}^{2}-\theta_{1}^{2}\right)\right) \\
H(a) & =\mathcal{L}\left(a\left(\theta_{0}-\theta_{1}\right),-\frac{1}{2} a\left(\theta_{0}^{2}-\theta_{1}^{2}\right)\right) \tag{5.2}
\end{align*}
$$

As in section 2, the function $L$ is steep. Therefore, we obtain an SLDP for $\left(V_{T}\right)$ similar to Theorem 3.1.

THEOREM 5.1. Under the hypothesis $H_{0},\left(T^{-1} V_{T}\right)$ satisfies an SLDP associated with $L$ and $H$. More precisely, for all $c>-\left(\theta_{0}-\theta_{1}\right)^{2} /\left(4 \theta_{0}\right)$, there exists a sequence $\left(b_{c, k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
\mathbf{P}\left\{V_{T} \geqq c T\right\}=\frac{e^{-T I(c)+H\left(a_{c}\right)}}{\sigma_{c} a_{c} \sqrt{2 \pi T}}\left[1+\sum_{k=1}^{p} \frac{b_{c, k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right] \tag{5.3}
\end{equation*}
$$

with $a_{c}$ given by $L^{\prime}\left(a_{c}\right)=c$ and

$$
\begin{equation*}
\sigma_{c}^{2}=L^{\prime \prime}\left(a_{c}\right)=-\frac{\left(\theta_{0}-\theta_{1}+2 c\right)^{3}}{\theta_{0}^{2}-\theta_{1}^{2}} \tag{5.4}
\end{equation*}
$$

The coefficients $b_{c, 1}, b_{c, 2}, \ldots, b_{c, p}$ may be explicitly given as in Theorem 3.1.

## 6. Proofs of the main results.

6.1. Proof of Theorem 3.1. Let $L_{T}$ be the normalized c.g.f. of $S_{T}$.

For all $c>-1 /(2 \theta)$, set $\mathbf{P}\left\{S_{T} \geqq c T\right\}=A_{T} B_{T}$ with

$$
\begin{align*}
A_{T} & =\exp \left[T\left(L_{T}\left(a_{c}\right)-c a_{c}\right)\right]  \tag{6.1}\\
B_{T} & =\mathbf{E}_{T}\left(\exp \left[-a_{c}\left(S_{T}-c T\right)\right] \mathbb{I}_{\left\{S_{T} \geqq c T\right\}}\right) \tag{6.2}
\end{align*}
$$

where $\mathbf{E}_{T}$ is the expectation after the usual change of probability

$$
\begin{equation*}
\frac{d \mathbf{Q}_{T}}{d \mathbf{P}}=\exp \left(a_{c} S_{T}-T L_{T}\left(a_{c}\right)\right) \tag{6.3}
\end{equation*}
$$

For all $a<\theta^{2} / 2$, set $R_{T}(a)=\mathcal{R}_{T}(0, a)$. It follows from part (ii) of Lemma 2.1 that $R_{T}\left(a_{c}\right)=\mathcal{O}\left(e^{-T /|c|}\right)$. In addition, we also have from (2.1) together with (3.2) that

$$
\begin{equation*}
A_{T}=\exp \left[-T I(c)+H\left(a_{c}\right)\right]\left(1+\mathcal{O}\left(e^{-T /|c|}\right)\right) \tag{6.4}
\end{equation*}
$$

It now remains to give an expansion of $B_{T}$ which can be rewritten as

$$
\begin{equation*}
B_{T}=\mathbf{E}_{T}\left(\exp \left[-a_{c} \sigma_{c} \sqrt{T} U_{T}\right] \mathbb{1}_{\left\{U_{T} \geqq 0\right\}}\right), \quad \text { where } U_{T}=\frac{S_{T}-c T}{\sigma_{c} \sqrt{T}} \tag{6.5}
\end{equation*}
$$

Lemma 6.1. For $c>-1 /(2 \theta)$, the distribution of $U_{T}$ under $\mathbf{Q}_{T}$ converges, as $T$ goes to infinity, to the $N(0,1)$ distribution. Moreover, there exists a sequence $\left(\delta_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
B_{T}=\frac{1}{\sqrt{T}}\left(\sum_{k=0}^{p} \frac{\delta_{k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right) \tag{6.6}
\end{equation*}
$$

The sequence $\left(\delta_{k}\right)$ depends only on the derivatives of $L$ and $H$ at point $a_{c}$. For example,

$$
\delta_{0}=\frac{1}{\sigma_{c} a_{c} \sqrt{2 \pi}} \quad \text { and } \quad \delta_{1}=\frac{-2 c \delta_{1}\left(12 \theta^{4} c^{4}+12 \theta^{3} c^{3}+35 \theta^{2} c^{2}+4 \theta c+2\right)}{\left(4 \theta^{2} c^{2}-1\right)^{2}}
$$

Proof of Theorem 3.1. Relation (3.3) follows from the conjunction of (6.4) and (6.6), which completes the proof of Theorem 3.1.
6.2. Proof of Theorem 4.1. This proof is completely different from the previous one since the function $L$ is steep for $\theta<c \leqq \theta / 2$, while this is no longer true when $c>\theta / 2$. The point $a_{c}$, given by (4.7), belongs to $\Gamma$ whenever $c<\theta / 3$. First, if $\theta<c \leqq \theta / 3$, we prove (4.6) as (3.3) via the usual change of probability.

Next, if $c>\theta / 3$, we use a slight modification of the strategy of time varying change of probability proposed in [10] and [8]. Let $L_{T}$ be the normalized c.g.f. of $Z_{T}(c)$, where the parameter $c$ is omitted in order to simplify the notation. There is a unique $a_{T}$, which belongs to $\Gamma$ and converges to $a^{c}$ as $T \rightarrow \infty$, solution of

$$
\begin{equation*}
L^{\prime}\left(a_{T}\right)+\frac{H^{\prime}\left(a_{T}\right)}{T}=0 \tag{6.7}
\end{equation*}
$$

Set

$$
\begin{equation*}
\frac{d \mathbf{Q}_{T}}{d \mathbf{P}}=\exp \left(a_{T} Z_{T}(c)-T L_{T}\left(a_{T}\right)\right) \tag{6.8}
\end{equation*}
$$

and denote by $\mathbf{E}_{T}$ the expectation under this new probability. We have the decomposition $\mathbf{P}\left\{\hat{\theta}_{T} \geqq c\right\}=A_{T} B_{T}$ with

$$
\begin{align*}
A_{T} & =\exp \left[T L_{T}\left(a_{T}\right)\right]  \tag{6.9}\\
B_{T} & =\mathbf{E}_{T}\left(\exp \left[-a_{T} Z_{T}(c)\right] \mathbb{I}_{\left\{Z_{T}(c) \geqq 0\right\}}\right) \tag{6.10}
\end{align*}
$$

It follows from the identity (6.7) together with (4.3) that

$$
\begin{equation*}
T\left(\rho\left(a_{T}\right)-\left(a_{T}+\theta\right)\right)=\frac{\theta^{2}-c \theta+a_{T} c}{\rho\left(a_{T}\right)\left(\rho\left(a_{T}\right)+c\right)} \tag{6.11}
\end{equation*}
$$

Since $\rho\left(a_{T}\right)=\sqrt{\theta^{2}+2 a_{T} c}$, we have $\lim _{T \rightarrow \infty} \rho\left(a_{T}\right)=2 c-\theta$, and (6.11) immediately implies

$$
\begin{equation*}
\lim _{T \rightarrow \infty} T\left(\rho\left(a_{T}\right)-\left(a_{T}+\theta\right)\right)=\frac{c-\theta}{3 c-\theta} \tag{6.12}
\end{equation*}
$$

From (6.7), there exists a sequence $\left(a_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
a_{T}=a^{c}+\sum_{k=1}^{p} \frac{a_{k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right) \tag{6.13}
\end{equation*}
$$

For example, by (6.12),

$$
a_{1}=-\frac{2 c-\theta}{3 c-\theta} \quad \text { and } \quad a_{2}=-\frac{c\left(c^{2}-5 \theta c+2 \theta^{2}\right)}{2(c-\theta)(3 c-\theta)^{3}}
$$

Moreover, for all $a \in \Gamma$, set $R_{T}(a)=\mathcal{R}_{T}(a,-c a)$. Using (4.3), relation (2.1) can be rewritten as

$$
\begin{equation*}
L_{T}(a)=L(a)+\frac{H(a)}{T}+\frac{1}{T} R_{T}(a) \tag{6.14}
\end{equation*}
$$

with $a \in \Gamma$. Next, by part (ii) of Lemma 2.1, the remainder $R_{T}\left(a_{T}\right)=\mathcal{O}\left(T e^{-2(2 c-\theta) T}\right)$. Thus, we obtain from (6.9) and (6.14) that

$$
\begin{equation*}
A_{T}=\exp \left[T L\left(a_{T}\right)+H\left(a_{T}\right)\right]\left(1+\mathcal{O}\left(T e^{-2(2 c-\theta) T}\right)\right) \tag{6.15}
\end{equation*}
$$

On the one hand, it follows from (6.13) that there exists a sequence $\left(\alpha_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
T L\left(a_{T}\right)=T L\left(a^{c}\right)+\frac{1}{2}+\sum_{k=1}^{p} \frac{\alpha_{k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right) \tag{6.16}
\end{equation*}
$$

The sequence $\left(\alpha_{k}\right)$ depends only on $\left(a_{k}\right)$ and on the derivatives of $L$ at point $a^{c}$. For example, $\alpha_{1}=a_{2} L^{\prime}\left(a^{c}\right)+\frac{1}{2} a_{1}^{2} L^{\prime \prime}\left(a^{c}\right)$ so that

$$
\alpha_{1}=\frac{c\left(2 c^{3}+c^{2}(1-3 \theta)+\theta c(\theta-5)+2 \theta^{2}\right)}{4(c-\theta)(2 c-\theta)(3 c-\theta)^{2}}
$$

On the other hand, we also have from (2.2) and (4.3) that

$$
\begin{equation*}
\exp \left[H\left(a_{T}\right)\right]=\sqrt{\frac{2 T \rho\left(a_{T}\right)}{T\left(\rho\left(a_{T}\right)-\left(a_{T}+\theta\right)\right)}} \tag{6.17}
\end{equation*}
$$

Hence, from (6.12) together with (6.13), there exists a sequence $\left(\beta_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
\exp \left[H\left(a_{T}\right)\right]=\sqrt{T} \sqrt{\frac{2(2 c-\theta)(3 c-\theta)}{c-\theta}}\left(1+\sum_{k=1}^{p} \frac{\beta_{k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right) . \tag{6.18}
\end{equation*}
$$

The sequence $\left(\beta_{k}\right)$ can be explicitly given as $\left(\alpha_{k}\right)$. For example,

$$
\beta_{1}=-\frac{c\left(c^{2}-3 \theta c+\theta^{2}\right)}{(c-\theta)(2 c-\theta)(3 c-\theta)^{2}}
$$

Finally, from (6.15) together with (6.16) and (6.18) it follows that there exists a sequence $\left(\gamma_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
A_{T}=\exp [-T I(c)] \sqrt{e T} \sqrt{\frac{2(2 c-\theta)(3 c-\theta)}{c-\theta}}\left(1+\sum_{k=1}^{p} \frac{\gamma_{k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right)\right) \tag{6.19}
\end{equation*}
$$

The sequence $\left(\gamma_{k}\right)$ depends only on $\left(a_{k}\right)$ and on the derivatives of $L$ at point $a^{c}$. For example, $\gamma_{1}=\alpha_{1}+\beta_{1}$ so that

$$
\gamma_{1}=\frac{c\left(2 c^{3}-3 c^{2}(1+\theta)+\theta c(\theta+7)-2 \theta^{2}\right)}{4(c-\theta)(2 c-\theta)(3 c-\theta)^{2}}
$$

For $c>\theta / 3$, the following lemma gives the asymptotical behavior of the distribution of $U_{T}$ and an expansion of $B_{T}$ which can be rewritten as

$$
\begin{equation*}
B_{T}=\mathbf{E}_{T}\left(\exp \left[-a_{T} T U_{T}\right] \mathbb{I}_{U_{T} \geqq 0}\right), \quad \text { where } U_{T}=\frac{Z_{T}(c)}{T} \tag{6.20}
\end{equation*}
$$

Lemma 6.2. (i) The distribution of $U_{T}$ under $\mathbf{Q}_{T}$ converges, as $T$ goes to infinity, to the distribution of $\gamma\left(N^{2}-1\right)$, where $N$ is an $N(0,1)$ random variable and $\gamma=$ $-L^{\prime}\left(a^{c}\right)=(3 c-\theta) /(2(2 c-\theta))$; i.e., the limit of the characteristic function of $U_{T}$ under $\mathbf{Q}_{T}$ is

$$
\begin{equation*}
\Phi(u)=\frac{\exp (-i \gamma u)}{\sqrt{1-2 i \gamma u}} \tag{6.21}
\end{equation*}
$$

(ii) There exists a sequence $\left(\delta_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
B_{T}=\sum_{k=1}^{p} \frac{\delta_{k}}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right) \tag{6.22}
\end{equation*}
$$

The sequence $\left(\delta_{k}\right)$ depends only on the Taylor expansion of $a_{T}$ at the neighborhood of $a^{c}$ together with the derivatives of $L$ at point $a^{c}$. For example,

$$
\delta_{1}=\frac{1}{a^{c} \gamma \sqrt{2 \pi e}} \quad \text { and } \quad \delta_{2}=\frac{\delta_{1}}{\gamma} \frac{-8 c^{3}+8 c^{2} \theta-5 c \theta^{2}+\theta^{3}}{2(c-\theta)(2 c-\theta)(3 c-\theta)^{2}}
$$

Remark. Part (i) was previously proven in [14] (see Lemma 4.6, p. 17). Part (ii) is not a consequence of (i). It requires a more precise study of the convergence of $\Phi_{T}$ towards $\Phi$ given in Lemma 7.2 below.

Proof of Theorem 4.1. Relation (4.8) follows from the conjunction of (6.19) and (6.22), which completes the proof of Theorem 4.1.
6.3. Proof of Theorem 4.2. We follow the same approach as that of Theorem 4.1. Assume that $c=\theta / 3$ so that $a_{c}=a^{c}=a_{\theta}$ with $a_{\theta}=-4 \theta / 3$. There is a unique $a_{T}$ solution, which belongs to $\Gamma$ and converges to $a_{\theta}$ as $T \rightarrow \infty$, of the equation

$$
\begin{equation*}
L^{\prime}\left(a_{T}\right)+\frac{H^{\prime}\left(a_{T}\right)}{T}=0 \tag{6.23}
\end{equation*}
$$

From (2.2) together with (4.3), it is easy to see that

$$
\begin{equation*}
\lim _{T \rightarrow \infty}\left(\rho\left(a_{T}\right)-\left(a_{T}+\theta\right)\right) H^{\prime}\left(a_{T}\right)=1 \tag{6.24}
\end{equation*}
$$

Moreover, a Taylor expansion around $a_{\theta}$ gives $L^{\prime}\left(a_{T}\right)=\left(a_{T}-a_{\theta}\right) \sigma_{\theta}^{2}+o\left(a_{T}-a_{\theta}\right)$ with $\sigma_{\theta}^{2}=L^{\prime \prime}\left(a_{\theta}\right)=-3 /(2 \theta)$. Thus, we find from (6.23) together with (6.24)

$$
\begin{equation*}
\lim _{T \rightarrow \infty} T\left(a_{T}-a_{\theta}\right)^{2}=\frac{1}{2 \sigma_{\theta}^{2}} \tag{6.25}
\end{equation*}
$$

Next, as in section 6.2, there exists a sequence $\left(a_{k}\right)$ such that, for any $p>0$ and $T$ large enough,

$$
\begin{equation*}
a_{T}=a_{\theta}+\sum_{k=1}^{2 p} \frac{a_{k}}{(\sqrt{T})^{k}}+\mathcal{O}\left(\frac{1}{T^{p} \sqrt{T}}\right) \tag{6.26}
\end{equation*}
$$

For example, $a_{1}=-1 /\left(\sigma_{\theta} \sqrt{2}\right)$ and $a_{2}=-\frac{1}{8}$. In addition, if we use the decomposition $\mathbf{P}\left\{\hat{\theta}_{T} \geqq c\right\}=A_{T} B_{T}$ given by (6.9) and (6.10), we obtain, as in (6.19), that for any $p>0$ and $T$ large enough

$$
\begin{equation*}
A_{T}=\exp [-T I(c)](e T)^{1 / 4}\left(-\frac{\theta}{3}\right)^{1 / 4}\left(1+\sum_{k=1}^{2 p} \frac{\gamma_{k}}{(\sqrt{T})^{k}}+\mathcal{O}\left(\frac{1}{T^{p} \sqrt{T}}\right)\right) \tag{6.27}
\end{equation*}
$$

where the sequence $\left(\gamma_{k}\right)$ may be explicitly calculated. For $c=\theta / 3$, it now remains to give an expansion of $B_{T}$ which can be rewritten as

$$
\begin{equation*}
B_{T}=\mathbf{E}_{T}\left(\exp \left[-a_{T} \sqrt{T} U_{T}\right] \mathbb{I}_{\left\{U_{T} \geqq 0\right\}}\right), \quad \text { where } U_{T}=\frac{Z_{T}(c)}{\sqrt{T}} \tag{6.28}
\end{equation*}
$$

Lemma 6.3. (i) The distribution of $U_{T}$ under $\mathbf{Q}_{T}$ converges, as $T$ goes to infinity, to the distribution of $\sigma_{\theta} N_{1}+\eta_{\theta}\left(N_{2}^{2}-1\right)$, where $N_{1}$ and $N_{2}$ are independent $N(0,1)$ random variables and $\sigma_{\theta}^{2}=L^{\prime \prime}\left(a_{\theta}\right)=-3 /(2 \theta)$ and $\eta_{\theta}=\sigma_{\theta} / \sqrt{2}$; i.e., the limit of the characteristic function of $U_{T}$ under $\mathbf{Q}_{T}$ is

$$
\begin{equation*}
\Phi(u)=\frac{\exp \left(-i \eta_{\theta} u-\sigma_{\theta}^{2} u^{2} / 2\right)}{\sqrt{1-2 i \eta_{\theta} u}} \tag{6.29}
\end{equation*}
$$

(ii) For any $p>0$ and $T$ large enough

$$
\begin{equation*}
B_{T}=\sum_{k=1}^{2 p} \frac{\delta_{k}}{(\sqrt{T})^{k}}+\mathcal{O}\left(\frac{1}{T^{p} \sqrt{T}}\right) \tag{6.30}
\end{equation*}
$$

where the sequence $\left(\delta_{k}\right)$ may be explicitly calculated. For example,

$$
\delta_{1}=\frac{1}{4 \pi a_{\theta} \eta_{\theta}} e^{-1 / 4} \Gamma\left(\frac{1}{4}\right)
$$

Proof of Theorem 4.2. Relation (4.12) follows from the conjunction of (6.27) and (6.30), which completes the proof of Theorem 4.2.
6.4. Proof of Theorem 5.1. The proof follows exactly along the same lines as that of Theorem 3.1 since the function $L$ associated with the log-likelihood ratio $\left(V_{T}\right)$ is steep.

## 7. Proofs of lemmas.

7.1. Proof of Lemma 2.2. (i) We refer to Janson [17, Chaps. 2 and 6]. The two random variables $Y_{1}=X_{T}^{2} / 2$ and $Y_{2}=\int_{0}^{T} X_{t}^{2} d t$ are in the chaos $H^{: 2:} \oplus H^{: 0}$. The operators $B_{1}$ and $B_{2}$ associated with their bilinear forms are positive, trace class, and their traces are expectations $\operatorname{tr} B_{1}=\mathbf{E} Y_{1}$ and $\operatorname{tr} B_{2}=\mathbf{E} Y_{2}$. By linearity, for any $(a, b) \in \mathbf{R}^{2}, a\left(Y_{1}-\mathbf{E} Y_{1}\right)+b\left(Y_{2}-\mathbf{E} Y_{2}\right)$ belongs to $H^{: 2}$. Consequently, from Theorem 6.1 of [17, p. 78], we have the decomposition

$$
\begin{equation*}
a\left(Y_{1}-\mathbf{E} Y_{1}\right)+b\left(Y_{2}-\mathbf{E} Y_{2}\right)=\sum_{j=1}^{\infty} \lambda_{j}^{T}\left(\varepsilon_{j}^{2}-1\right) \tag{7.1}
\end{equation*}
$$

where $\left(\varepsilon_{j}\right)$ are independent $N(0,1)$ random variables and $\left(\lambda_{j}^{T}\right)$ are the eigenvalues of $a B_{1}+b B_{2}$. Since this operator is trace class, part (i) of Lemma 2.2 follows from (7.1).
(ii) The spectrum $\sigma\left(B_{2}\right) \subset\left[0, \theta^{-2}\right]$ (see [8]) and $\mathbf{E} X_{T}^{2} / 2$ is uniformly bounded in $T$.
(iii) From (1.6), (1.7), and (2.5), the normalized c.g.f. of $\mathcal{Z}_{T}(a, b)$ in $\alpha$ is $\mathcal{L}_{T}(\alpha a, \alpha b)$. From part (i), it can be rewritten as

$$
\begin{equation*}
\mathcal{L}_{T}(\alpha a, \alpha b)=-\frac{\alpha a}{2}-\frac{1}{2}\left\langle\nu_{T}, \Psi_{\alpha}\right\rangle, \tag{7.2}
\end{equation*}
$$

where $\nu_{T}$ is given by $(2.6)$ and $\Psi_{\alpha}(x)=\log (1-2 \alpha x)$ (notice that $\left.\Psi_{\alpha} \in \mathcal{F}\right)$. From Lemma 2.1, we already know that

$$
\begin{equation*}
\lim _{T \rightarrow \infty} \mathcal{L}_{T}(\alpha a, \alpha b)=\mathcal{L}(\alpha a, \alpha b)=-\frac{1}{2}(\alpha a+\theta+\rho(\alpha b)) \tag{7.3}
\end{equation*}
$$

for $\alpha$ in some interval $\left[-\alpha_{0}, \alpha_{0}\right]$ depending on $a$ and $b$, with $\alpha_{0}>0$. Thus, it immediately follows from (7.2) and (7.3) that $\lim _{T \rightarrow \infty}\left\langle\nu_{T}, \Psi_{\alpha}\right\rangle=\theta+\rho(\alpha b)$. Since $\nu$ is defined by $\langle\nu, h\rangle=1 /(2 \pi) \int_{\mathbf{R}} h(b g(x)) d x$, where $h$ is any continuous function with compact support, an easy integration by parts gives $\left\langle\nu, \Psi_{\alpha}\right\rangle=\theta+\rho(\alpha b)$ which proves that

$$
\begin{equation*}
\lim _{T \rightarrow \infty}\left\langle\nu_{T}, \Psi_{\alpha}\right\rangle=\left\langle\nu, \Psi_{\alpha}\right\rangle \tag{7.4}
\end{equation*}
$$

for every $\alpha \in\left[-\alpha_{0}, \alpha_{0}\right]$. By use of the above part (ii) together with a slight modification of Lemma 9 of [4], we arrive from (7.4) at the convergence of $\nu_{T}$ towards $\nu$, which completes the proof of Lemma 2.2.
7.2. Proof of Lemma 6.1. If $\Phi_{T}$ is the characteristic function of $U_{T}$ under $\mathbf{Q}_{T}$, (6.3) immediately implies

$$
\begin{equation*}
\Phi_{T}(u)=\exp \left[-\frac{i u \sqrt{T} c}{\sigma_{c}}+T\left(L_{T}\left(a_{c}+\frac{i u}{\sigma_{c} \sqrt{T}}\right)-L_{T}\left(a_{c}\right)\right)\right] \tag{7.5}
\end{equation*}
$$

so that

$$
\begin{equation*}
\left|\Phi_{T}(u)\right|^{2}=\prod_{j=1}^{\infty}\left(1+\frac{4 u^{2}\left(\lambda_{j}^{T}\right)^{2}}{\sigma_{c}^{2} T\left(1-2 a_{c} \lambda_{j}^{T}\right)^{2}}\right)^{-1 / 2} \tag{7.6}
\end{equation*}
$$

Choose $\varepsilon>0$ such that $1-2 a_{c} \varepsilon>0$ and let $q_{T}=\operatorname{card}\left\{\lambda_{j}^{T} \mid \lambda_{j}^{T}>\varepsilon\right\}$. Then, by Lemma 2.2, part (iii), there exists some positive constant $\eta$, depending only on $\varepsilon$, such that

$$
\begin{equation*}
\liminf _{T \rightarrow \infty} \frac{q_{T}}{T} \geqq \eta \tag{7.7}
\end{equation*}
$$

Thus, it follows from (7.6) together with (7.7) that, for $T$ large enough,

$$
\begin{equation*}
\left|\Phi_{T}(u)\right|^{2} \leqq\left(1+\frac{\xi u^{2}}{T}\right)^{-\eta T / 2} \tag{7.8}
\end{equation*}
$$

with $\xi>0$. Thereby, for $T$ large enough, $\Phi_{T} \in L^{2}(\mathbf{R})$ and we can use the Parseval formula in (6.5) to get

$$
\begin{equation*}
B_{T}=\frac{1}{2 \pi a_{c} \sigma_{c} \sqrt{T}} \int_{\mathbf{R}}\left(1+\frac{i u}{a_{c} \sigma_{c} \sqrt{T}}\right)^{-1} \Phi_{T}(u) d u \tag{7.9}
\end{equation*}
$$

For some positive constant $s$, set $s_{T}=s T^{1 / 6}$. We separate $B_{T}=C_{T}+D_{T}$, where

$$
\begin{align*}
C_{T} & =\frac{1}{2 \pi a_{c} \sigma_{c} \sqrt{T}} \int_{|u| \leqq s_{T}}\left(1+\frac{i u}{a_{c} \sigma_{c} \sqrt{T}}\right)^{-1} \Phi_{T}(u) d u  \tag{7.10}\\
D_{T} & =\frac{1}{2 \pi a_{c} \sigma_{c} \sqrt{T}} \int_{|u|>s_{T}}\left(1+\frac{i u}{a_{c} \sigma_{c} \sqrt{T}}\right)^{-1} \Phi_{T}(u) d u \tag{7.11}
\end{align*}
$$

On the one hand, we find from (7.8) that

$$
\begin{equation*}
\left|D_{T}\right|=O\left(\left(1+\frac{\xi s_{T}^{2}}{T}\right)^{-\eta T / 4}\right) \tag{7.12}
\end{equation*}
$$

so that $\left|D_{T}\right|=\mathcal{O}\left(e^{-\mu T^{1 / 3}}\right)$ with $\mu>0$. On the other hand, it follows from (2.3) that for any $k \in \mathbf{N}, R_{T}^{(k)}\left(a_{c}\right)=\mathcal{O}\left(T^{k} e^{-T /|c|}\right)$. Then, using (3.2) and (2.1) we get

$$
\begin{equation*}
L_{T}^{(k)}\left(a_{c}\right)=L^{(k)}\left(a_{c}\right)+\frac{H^{(k)}\left(a_{c}\right)}{T}+\mathcal{O}\left(T^{k} e^{-T /|c|}\right) \tag{7.13}
\end{equation*}
$$

Consequently, via (7.5) together with (7.13) we can prove the following Taylor expansion for $\Phi_{T}$. Similar expansion was established in the i.i.d. case by Cramer [7] and Esseen [12].

Lemma 7.1. For any $p>0$, there exist integers $q(p), r(p)$, and a sequence $\left(\varphi_{k, l}\right)$ independent of $p$ such that, for $T$ large enough,

$$
\begin{equation*}
\Phi_{T}(u)=\exp \left(-\frac{u^{2}}{2}\right)\left[1+\frac{1}{\sqrt{T}} \sum_{k=0}^{2 p} \sum_{l=k+1}^{q(p)} \frac{\varphi_{k, l} u^{l}}{(\sqrt{T})^{k}}+\mathcal{O}\left(\frac{\max \left(1,|u|^{r(p)}\right)}{T^{p+1}}\right)\right] \tag{7.14}
\end{equation*}
$$

and the remainder $\mathcal{O}$ is uniform as soon as $|u|=\mathcal{O}\left(T^{1 / 6}\right)$.
Proof. From (7.5) and (7.13), we have

$$
\begin{aligned}
\log \Phi_{T}(u)= & -\frac{u^{2}}{2}+T \sum_{k=3}^{2 p+3}\left(\frac{i u}{\sigma_{c} \sqrt{T}}\right)^{k} \frac{L^{(k)}\left(a_{c}\right)}{k!}+\sum_{k=1}^{2 p+1}\left(\frac{i u}{\sigma_{c} \sqrt{T}}\right)^{k} \frac{H^{(k)}\left(a_{c}\right)}{k!} \\
& +\mathcal{O}\left(\frac{\max \left(1, u^{2 p+4}\right)}{T^{p+1}}\right)
\end{aligned}
$$

Thus, we follow the same approach as Cramer (see [7, Lemma 2, p. 72]), remarking that in the range $|u|=\mathcal{O}\left(T^{1 / 6}\right)$ the quantity $u^{l} /(\sqrt{T})^{k}$ remains bounded in (7.14). Lemma 7.1 is proved.

Proof of Lemma 6.1. Relation (6.6) follows from (7.10) and (7.14) together with standard calculus on the $N(0,1)$ distribution.
7.3. Proof of Lemma 6.2. Let $\Phi_{T}$ be the characteristic function of $U_{T}$ un$\operatorname{der} \mathbf{Q}_{T}$. We have from (6.8)

$$
\begin{equation*}
\Phi_{T}(u)=\exp \left[T\left(L_{T}\left(a_{T}+\frac{i u}{T}\right)-L_{T}\left(a_{T}\right)\right)\right] \tag{7.16}
\end{equation*}
$$

so that

$$
\begin{equation*}
\left|\Phi_{T}(u)\right|^{2}=\prod_{j=1}^{\infty}\left(1+\frac{4 u^{2}\left(\lambda_{j}^{T}\right)^{2}}{T^{2}\left(1-2 a_{T} \lambda_{j}^{T}\right)^{2}}\right)^{-1 / 2} \tag{7.17}
\end{equation*}
$$

As in (7.8), it follows from Lemma 2.2, part (iii) that, for some positive constants $\eta, \xi$, and for $T$ large enough,

$$
\begin{equation*}
\left|\Phi_{T}(u)\right|^{2} \leqq\left(1+\frac{\xi u^{2}}{T^{2}}\right)^{-\eta T} \tag{7.18}
\end{equation*}
$$

Thereby, for $T$ large enough, $\Phi_{T} \in L^{2}(\mathbf{R})$ and we can use the Parseval formula in (6.20) to obtain that

$$
\begin{equation*}
B_{T}=\frac{1}{2 \pi T a_{T}} \int_{\mathbf{R}}\left(1+\frac{i u}{T a_{T}}\right)^{-1} \Phi_{T}(u) d u \tag{7.19}
\end{equation*}
$$

Let $s_{T}>0$ be such that $\sqrt{T}=o\left(s_{T}\right)$ as $T \rightarrow \infty$. We split $B_{T}$ into two terms, $B_{T}=C_{T}+D_{T}$, where

$$
\begin{align*}
C_{T} & =\frac{1}{2 \pi T a_{T}} \int_{|u| \leqq s_{T}}\left(1+\frac{i u}{T a_{T}}\right)^{-1} \Phi_{T}(u) d u  \tag{7.20}\\
D_{T} & =\frac{1}{2 \pi T a_{T}} \int_{|u|>s_{T}}\left(1+\frac{i u}{T a_{T}}\right)^{-1} \Phi_{T}(u) d u \tag{7.21}
\end{align*}
$$

First, we find as in (7.12) that, if $s_{T}=o(T),\left|D_{T}\right|=\mathcal{O}\left(\exp \left(-\mu s_{T}^{2} / T\right)\right)$ with $\mu>0$. It now remains to precisely evaluate $C_{T}$ via the following Taylor expansion for $\Phi_{T}$.

LEMMA 7.2. For any $p>0$, there exist integers $q(p), r(p), s(p)$, and a sequence $\left(\varphi_{k, l, m}\right)$ independent of $p$ such that, for $T$ large enough,

$$
\begin{aligned}
\Phi_{T}(u)= & \Phi(u) \exp \left(-\frac{\sigma^{2} u^{2}}{2 T}\right) \\
& \times\left[1+\sum_{k=1}^{p} \sum_{l=k+1}^{q(p)} \sum_{m=0}^{r(p)} \frac{\varphi_{k, l, m} u^{l}}{T^{k}(1-2 i \gamma u)^{m}}+\mathcal{O}\left(\frac{\max \left(1,|u|^{s(p)}\right)}{T^{p+1}}\right)\right]
\end{aligned}
$$

where $\Phi$ is given by (6.21),

$$
\gamma=-L^{\prime}\left(a^{c}\right)=\frac{3 c-\theta}{2(2 c-\theta)} \quad \text { and } \quad \sigma^{2}=L^{\prime \prime}\left(a^{c}\right)=\frac{c^{2}}{2(2 c-\theta)^{3}}
$$

Moreover, the remainder $\mathcal{O}$ is uniform as soon as $|u|=o\left(T^{2 / 3}\right)$.
Remark. One can see in this asymptotic expansion the limit $\chi^{2}$ distribution $\Phi$ together with an independent centered Gaussian distribution with small variance $\sigma^{2} / T$.

Proof. From (6.14) together with (7.16)

$$
\begin{aligned}
\Phi_{T}(u)= & \exp \left[T\left(L\left(a_{T}+\frac{i u}{T}\right)-L\left(a_{T}\right)\right)\right] \\
& \times \exp \left[H\left(a_{T}+\frac{i u}{T}\right)-H\left(a_{T}\right)\right]\left(1+\mathcal{O}\left(e^{-\xi T}\right)\right)
\end{aligned}
$$

with $\xi>0$. On the one hand, (4.5) immediately implies

$$
\begin{equation*}
T\left(L\left(a_{T}+\frac{i u}{T}\right)-L\left(a_{T}\right)\right)=-\frac{i u}{2}-\frac{T}{2} \rho\left(a_{T}\right)\left(\left(1+\frac{i u b_{T}}{T}\right)^{1 / 2}-1\right) \tag{7.22}
\end{equation*}
$$

where $b_{T}=2 c\left(\rho\left(a_{T}\right)\right)^{-2}$. Consequently, for any $p \geqq 2$

$$
\begin{align*}
& \exp \left[T\left(L\left(a_{T}+\frac{i u}{T}\right)-L\left(a_{T}\right)\right)\right]=\exp \left(-\frac{i u}{4}\left(2+\rho\left(a_{T}\right) b_{T}\right)\right) \\
& \quad \times \exp \left(-\frac{T}{2} \rho\left(a_{T}\right) \sum_{k=2}^{p} l_{k}\left(\frac{i u b_{T}}{T}\right)^{k}\right)\left(1+\mathcal{O}\left(\frac{|u|^{p+1}}{T^{p+1}}\right)\right) \tag{7.23}
\end{align*}
$$

where $l_{k}=(-1)^{k-1}(2 k)!/\left((2 k-1)\left(2^{k} k!\right)^{2}\right)$. On the other hand, we obtain from (2.2)

$$
\begin{align*}
& \exp \left[H\left(a_{T}+\frac{i u}{T}\right)-H\left(a_{T}\right)\right] \\
& \quad=\sqrt{\frac{\rho\left(a_{T}\right)-\left(a_{T}+\theta\right)}{\rho\left(a_{T}\right)-\left(a_{T}+i u / T+\theta\right)\left(1+i u b_{T} / T\right)^{-1 / 2}}} . \tag{7.24}
\end{align*}
$$

If $c_{T}=T\left(\rho\left(a_{T}\right)-\left(a_{T}+\theta\right)\right)$ and $d_{T}(u)=1-i u / c_{T}+\left(a_{T}+\theta\right) i u b_{T} /\left(2 c_{T}\right)$, we have for any $p \geqq 2$

$$
\begin{align*}
& \exp \left[H\left(a_{T}+\frac{i u}{T}\right)\right.\left.-H\left(a_{T}\right)\right] \\
&=\frac{1}{\sqrt{d_{T}(u)}}\left(1-\frac{u^{2} b_{T}}{2 T c_{T} d_{T}(u)}-\frac{T\left(a_{T}+\theta+i u T^{-1}\right)}{c_{T} d_{T}(u)}\right. \\
&\left.\times\left[\sum_{k=2}^{p} h_{k}\left(\frac{i u b_{T}}{T}\right)^{k}+\mathcal{O}\left(\frac{|u|^{p+1}}{T^{p+1}}\right)\right]\right)^{-1 / 2} \tag{7.25}
\end{align*}
$$

with $h_{k}=(-1)^{k}(2 k)!/\left(2^{k} k!\right)^{2}$. As $T$ goes to infinity, the limits of $a_{T}, b_{T}, c_{T}$, and $d_{T}(u)$ are $a^{c}, 2 c /(2 c-\theta)^{2},(c-\theta) /(3 c-\theta)$, and $1-2 i \gamma u$, respectively. Therefore, we find by (7.23), together with (7.25), the pointwise convergence

$$
\begin{equation*}
\lim _{T \rightarrow \infty} \Phi_{T}(u)=\Phi(u)=\frac{\exp (-i \gamma u)}{\sqrt{1-2 i \gamma u}} \tag{7.26}
\end{equation*}
$$

which achieves the proof of the first part of Lemma 6.2. Finally, after tedious calculus, we prove Lemma 7.2 via a Taylor expansion of the exponential in (7.23) together with a Taylor expansion of the square root in (7.25).

End of the proof of Lemma 6.2. From (7.20) and the above expansion of $\Phi_{T}$, we have for $T$ large enough

$$
\begin{align*}
2 \pi T a_{T} C_{T}= & \int_{|u| \leqq s_{T}} \Phi(u) \exp \left(-\frac{\sigma^{2} u^{2}}{2 T}\right)\left[1+\mathcal{O}\left(\frac{\max \left(1,|u|^{s(p)}\right)}{T^{p+1}}\right)\right] d u \\
& +\sum_{k=1}^{p} \sum_{l=k+1}^{q(p)} \sum_{m=0}^{r(p)} \psi_{k, l, m} \int_{|u| \leqq s_{T}} \Phi(u) \exp \left(-\frac{\sigma^{2} u^{2}}{2 T}\right) \frac{u^{l}}{T^{k}(1-2 i \gamma u)^{m}} d u \tag{7.27}
\end{align*}
$$

where the sequence $\left(\varphi_{k, l, m}\right)$ is replaced by $\left(\psi_{k, l, m}\right)$ due to the factor preceding $\Phi_{T}$ in (7.20). Furthermore, for all $u \in \mathbf{R},|\Phi(u)| \leqq 1$ and

$$
\int_{|u|>s_{T}} \exp \left(-\frac{\sigma^{2} u^{2}}{2 T}\right) d u=\mathcal{O}\left(\exp \left(-\frac{\sigma^{2} s_{T}^{2}}{2 T}\right)\right)
$$

Thus, we may change in (7.27) all the integrals by integrals over $\mathbf{R}$ at the cost of $\mathcal{O}\left(\exp \left(-\mu s_{T}^{2} / T\right)\right)$ with $\mu>0$. We complete the proof of (6.22) via the following lemma which gives the values of integrals of the type

$$
\int_{\mathbf{R}} \exp \left(-i u \gamma-\frac{\sigma^{2} u^{2}}{2 T}\right) \frac{u^{\beta}}{(1-2 i \gamma u)^{\alpha}} d u
$$

Let $f_{\alpha}$ be the density of the gamma distribution with parameters $\alpha$ and $\frac{1}{2}$,

$$
f_{\alpha}(x)= \begin{cases}\exp \left(-\frac{x}{2}\right) \frac{x^{\alpha-1}}{2^{\alpha} \Gamma(\alpha)} & \text { if } x>0  \tag{7.28}\\ 0 & \text { otherwise }\end{cases}
$$

For any $p>0$ and $0 \leqq k \leqq p$, set

$$
\begin{equation*}
v_{k}(\alpha, \beta)=\frac{2 \pi \sigma^{2 k} i^{\beta}}{\gamma^{2 k+\beta+1} 2^{k} k!} f_{\alpha}^{(2 k+\beta)}(1) \tag{7.29}
\end{equation*}
$$

Lemma 7.3. For any $p>0$, we have

$$
\begin{equation*}
\int_{\mathbf{R}} \exp \left(-i u \gamma-\frac{\sigma^{2} u^{2}}{2 T}\right) \frac{u^{\beta}}{(1-2 i \gamma u)^{\alpha}} d u=\sum_{k=0}^{p} \frac{v_{k}(\alpha, \beta)}{T^{k}}+\mathcal{O}\left(\frac{1}{T^{p+1}}\right) \tag{7.30}
\end{equation*}
$$

Proof. Let $\widehat{f}_{\alpha}$ be the characteristic function of $f_{\alpha}, \widehat{f}_{\alpha}(u)=1 /(1-2 i u)^{\alpha}$. Denote by $N_{\tau}$ the Gaussian kernel of variance $\tau$. First, we notice that

$$
\begin{equation*}
\int_{\mathbf{R}} \exp \left(-i v \xi-\frac{\tau v^{2}}{2}\right) v^{\beta} \widehat{f}_{\alpha}(v) d v=2 \pi i^{\beta} f_{\alpha} * N_{\tau}^{(\beta)}(\xi) \tag{7.31}
\end{equation*}
$$

It is rather easy to see that, for every $p>0$,

$$
\begin{equation*}
f_{\alpha} * N_{\tau}^{(\beta)}(\xi)=\sum_{k=0}^{p} \frac{\tau^{k}}{2^{k} k!} f_{\alpha}^{(2 k+\beta)}(\xi)+\mathcal{O}\left(\tau^{(p+1)}\right) \tag{7.32}
\end{equation*}
$$

For the sake of completeness, let us prove (7.32). We start with a localization

$$
\left|f_{\alpha} * N_{\tau}^{(\beta)}(\xi)-\int_{|x| \leqq \delta} f_{\alpha}(\xi-x) N_{\tau}^{(\beta)}(x) d x\right| \leqq \sup _{|x|>\delta}\left|N_{\tau}^{(\beta)}(x)\right|
$$

with $0<\delta<\xi$. First, this supremum is $\mathcal{O}\left(\exp \left(-\mu_{1} \tau^{-1}\right)\right)$ with $\mu_{1}>0$. Next, the first term is easily integrated by parts since

$$
\int_{|x| \leqq \delta} f_{\alpha}(\xi-x) N_{\tau}^{(\beta)}(x) d x=\int_{|x| \leqq \delta} f_{\alpha}^{(\beta)}(\xi-x) N_{\tau}(x) d x+\mathcal{O}\left(\exp \left(-\mu_{2} \tau^{-1}\right)\right)
$$

with $\mu_{2}>0$. Furthermore, as the function $f_{\alpha} \in \mathcal{C}^{\infty}$ in $[\xi-\delta, \xi+\delta]$, we have the Taylor expansion

$$
f_{\alpha}^{(\beta)}(\xi-x)=\sum_{k=0}^{2 p+1} \frac{(-x)^{k}}{k!} f_{\alpha}^{(\beta+k)}(\xi)+\mathcal{O}\left(x^{2 p+2}\right)
$$

for every $x \in[-\delta,+\delta]$. This yields

$$
\begin{aligned}
\int_{|x| \leqq \delta} f_{\alpha}(\xi-x) N_{\tau}^{(\beta)}(x) d x= & \sum_{k=0}^{p} f_{\alpha}^{(2 k+\beta)}(\xi) \frac{1}{(2 k)!} \int_{|x| \leqq \delta} x^{2 k} N_{\tau}(x) d x \\
& +\mathcal{O}\left(\int_{|x| \leqq \delta} x^{2 p+2} N_{\tau}(x) d x\right)
\end{aligned}
$$

Therefore, if we remove the localization whose cost is $\mathcal{O}\left(\exp \left(-\mu_{3} \tau^{-1}\right)\right)$, we find that

$$
f_{\alpha} * N_{\tau}^{(\beta)}(\xi)=\sum_{k=0}^{p} f_{\alpha}^{(2 k+\beta)}(\xi) \frac{1}{(2 k)!} \int_{\mathbf{R}} x^{2 k} N_{\tau}(x) d x+\mathcal{O}\left(\int_{\mathbf{R}} x^{2 p+2} N_{\tau}(x) d x\right)
$$

which immediately leads to (7.32). Finally, by a change of variables in (7.31), we obtain

$$
\begin{equation*}
\int_{\mathbf{R}} \exp \left(-i u \gamma-\frac{\sigma^{2} u^{2}}{2 T}\right) \frac{u^{\beta}}{(1-2 i \gamma u)^{\alpha}} d u=\frac{2 \pi(-i)^{\beta}}{\gamma^{\beta+1}} f_{\alpha} * N_{\tau}^{(\beta)}(1) \tag{7.33}
\end{equation*}
$$

with $\tau=\sigma^{2} /\left(T \gamma^{2}\right)$, which completes the proof of Lemma 7.3.
7.4. Proof of Lemma 6.3. Let $\Phi_{T}$ be the characteristic function of $U_{T}$ un$\operatorname{der} \mathbf{Q}_{T}$ given by (6.8),

$$
\begin{equation*}
\Phi_{T}(u)=\exp \left[T\left(L_{T}\left(a_{T}+\frac{i u}{\sqrt{T}}\right)-L_{T}\left(a_{T}\right)\right)\right] . \tag{7.34}
\end{equation*}
$$

We prove only the pointwise convergence of $\Phi_{T}$ given in (6.29), as the proof of the Taylor expansion (6.30) follows essentially the same arguments as those of section 7.3 . From the definition (4.5) of $L$, we have

$$
\begin{equation*}
\lim _{T \rightarrow \infty} T\left(L\left(a_{T}+\frac{i u}{\sqrt{T}}\right)-L\left(a_{T}\right)\right)=-i \eta_{\theta} u-\frac{\sigma_{\theta}^{2} u^{2}}{2} . \tag{7.35}
\end{equation*}
$$

Moreover, we also have from the definition (2.2) of $H$

$$
\begin{equation*}
\lim _{T \rightarrow \infty} \exp \left[H\left(a_{T}+\frac{i u}{\sqrt{T}}\right)-H\left(a_{T}\right)\right]=\frac{1}{\sqrt{1-2 i \eta_{\theta} u}} \tag{7.36}
\end{equation*}
$$

Therefore, the pointwise convergence (6.29) immediately follows from (2.1), (7.35), and (7.36). The first term in (6.30) is $\delta_{1}=\left(2 \pi a_{\theta}\right)^{-1} \int_{\mathbf{R}} \Phi(u) d u$. By a change of variables together with the fact that $2 \eta_{\theta}^{2}=\sigma_{\theta}^{2}$, we find that

$$
\delta_{1}=\frac{1}{2 \pi a_{\theta} \eta_{\theta}} \int_{\mathbf{R}} \frac{1}{\sqrt{1-2 i u}} \exp \left(-\left(u^{2}+i u\right)\right) d u
$$

Finally, via a contour integral for the Gamma function, we obtain

$$
\delta_{1}=\frac{1}{4 \pi a_{\theta} \eta_{\theta}} \exp \left(-\frac{1}{4}\right) \Gamma\left(\frac{1}{4}\right)
$$

which completes the proof of Lemma 6.3.

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